

Monodromy in the Kepler Problem

Holger Dullin, University of Sydney,
School of Mathematics and Statistics

Geometric aspects of momentum maps and integrability
CSF Ascona, April 2018

joint work with Holger Waalkens, Groningen

HRD, Waalkens, PRL 120:020507 (2018)

Plan

- Review of Kepler problem
- (super) integrable systems
- Quantum integrable systems, joint spectrum
- Quantum monodromy in Kepler
- Toric and semi-toric systems
- Spheroidal harmonics (Laplacian)



Kepler problem (in 3D)

$$\mathbf{r} = (x, y, z)^t, \mathbf{p} = (\dot{x}, \dot{y}, \dot{z})^t, r = |\mathbf{r}|, \quad \dot{\mathbf{r}} = \mathbf{p}, \dot{\mathbf{p}} = -\frac{\mathbf{r}}{r^3}$$

- Liouville integrable: $H = \frac{1}{2}\mathbf{p}^2 - \frac{1}{r}, \quad \mathbf{L} = \mathbf{r} \times \mathbf{p}$
($H, |\mathbf{L}|^2, L_z$) three integrals in involution



Kepler problem (in 3D)

$$\mathbf{r} = (x, y, z)^t, \mathbf{p} = (\dot{x}, \dot{y}, \dot{z})^t, r = |\mathbf{r}|, \quad \dot{\mathbf{r}} = \mathbf{p}, \dot{\mathbf{p}} = -\frac{\mathbf{r}}{r^3}$$

- Liouville integrable: $H = \frac{1}{2}\mathbf{p}^2 - \frac{1}{r}, \quad \mathbf{L} = \mathbf{r} \times \mathbf{p}$
($H, |\mathbf{L}|^2, L_z$) three integrals in involution

- more integrals (super-integrable): $\mathbf{e} = \mathbf{p} \times \mathbf{L} - \frac{\mathbf{r}}{r}$
Runge-Lenz vector

- 7 integrals? Only 5 independent (in 6D)
 \Rightarrow EVERY orbit is periodic

- very special, but very important; similarly:
isotropic harmonic oscillator, free particle,
geodesic flow on spheres



Hydrogen atom

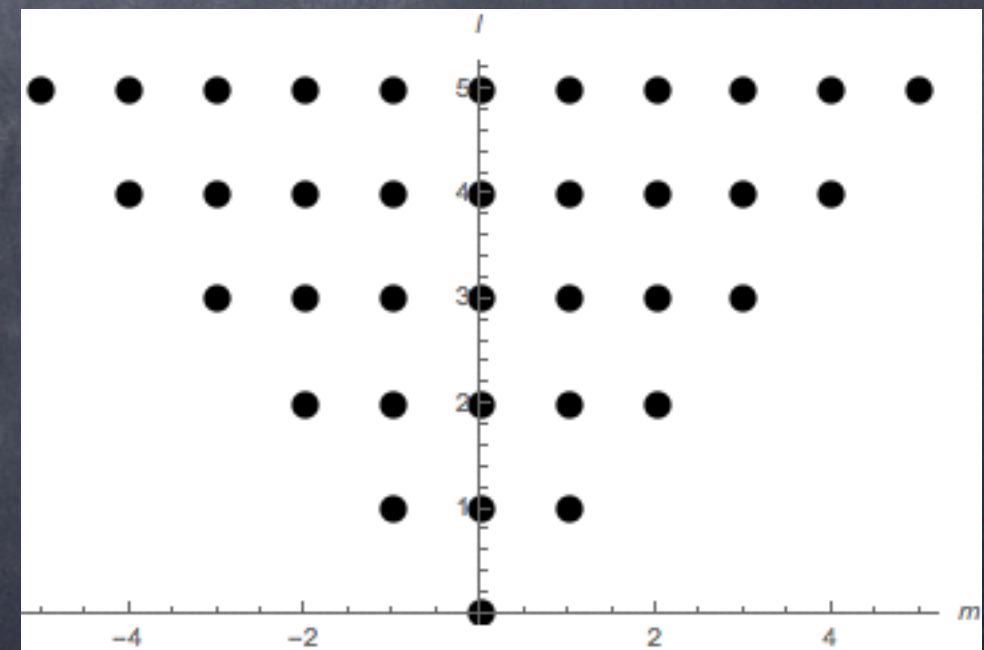
$$\mathbf{r} = (x, y, z)^t, \mathbf{p} = (\dot{x}, \dot{y}, \dot{z})^t, r = |\mathbf{r}|, \quad \dot{\mathbf{r}} = \mathbf{p}, \dot{\mathbf{p}} = -\frac{\mathbf{r}}{r^3}$$

- $H = \frac{1}{2}p^2 - \frac{1}{r}$ (Kepler in 3D)
- $\hat{H} = -\frac{1}{2}\hbar^2\nabla^2 - \frac{1}{r}$ (Schrödinger operator)
- $\hat{H}\psi = E\psi$ (Schrödinger equation, PDE)
- eigenvalues $E = -1/(2n^2)$, $n=1,2,3,\dots$
- degenerate eigenspace, multiplicity n^2



Quantum Integrable System

- quantum integrable system with classical limit
- separation of variables in spherical coordinates gives 3 commuting 2nd order differential operators $(H, |L|^2, L_z)$ on $L^2(\mathbb{R}^3)$
- joint spectrum of $(H, |L|^2, L_z)$ is lattice $n=6$:
 $(-1/(2n^2), l(l+1), m)$,
 $n=1,2,\dots, l=0,1,\dots,n-1, m=-l,\dots,l$
- for fixed n (e.g. $n=6$):



Hydrogen Orbitals

- separation of variables in spherical coordinates

- Laguerre polynomials for radial part

$n=4:$

$m=0$

1

2

3

- spherical harmonics $Y_l^m()$

- wave function Ψ_{nlm}

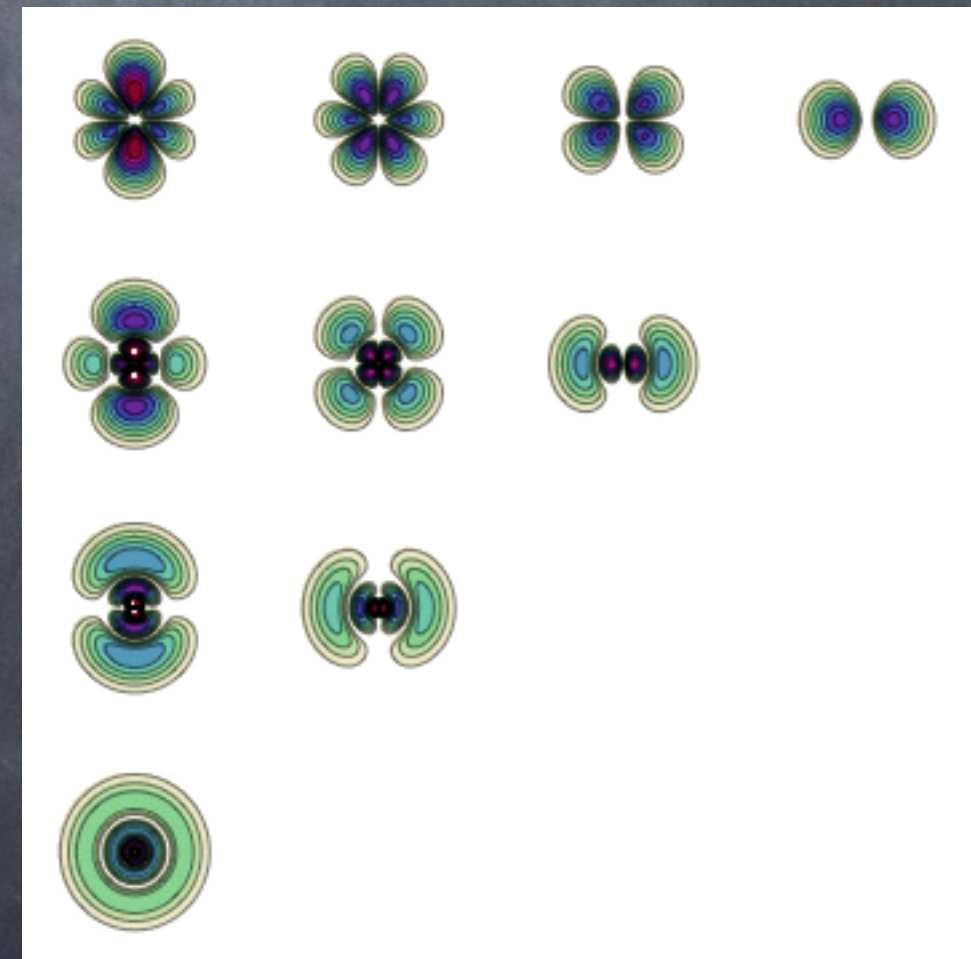
$l=3, f$

- slice of $|\Psi_{nlm}|$:

$l=2, d$

$l=1, p$

$l=0, s$



Super-integrable Kepler

- integrals: $H = \frac{1}{2}\mathbf{p}^2 - \frac{1}{r}, \quad \mathbf{L} = \mathbf{r} \times \mathbf{p}$
- more integrals: $\mathbf{e} = \mathbf{p} \times \mathbf{L} - \mathbf{r}/r$
Laplace-Runge-Lenz vector
- full symmetry algebra $\mathfrak{so}(4)$: $\mathbf{L}, \mathbf{K} = n\mathbf{e}, \quad H = -\frac{1}{2n^2}$
- Casimirs $\mathbf{L} \cdot \mathbf{K} = 0, \quad \mathbf{L}^2 + \mathbf{K}^2 = n^2$
- maximally super-integrable
every orbit is periodic
Integrals do not all have vanishing Poisson bracket!



$\mathfrak{so}(4)$ algebra

- either with $\{, \}$ (or in QM similarly with $[,]$):
- $\{ L_i, L_j \} = \epsilon_{ijk} L_k$
 $\{ L_i, K_j \} = \epsilon_{ijk} K_k$
 $\{ K_i, K_j \} = \epsilon_{ijk} L_k$
- Casimirs $(K+L)^2 = n^2$ and $(K-L)^2 = n^2$
- hence symplectic leaf is $S^2 \times S^2$
- a compact symplectic manifold
NB: reduction of maximally superintegrable systems with respect to Hamiltonian flow typically leads to compact symplectic mflds.



Super- vs. Liouville-integrable

- Liouville-integrable: n independent commuting integrals
- Conjecture that non-commutative integrability implies Liouville integrability (Mischenko-Fomenko '78)
- proved by Sadetov '03, Bolsinov '05
(but not proven in the same functional class)
- key to our analysis is this simple observation:

A superintegrable system may be Liouville integrable in non-equivalent ways



Liouville integrable realisations

- Definition: Given a super-integrable Hamiltonian H , a Liouville integrable realisation of H is a collection of functions (H, F_2, \dots, F_n) in involution, independent almost everywhere.
- For maximally super-integrable H there are infinitely many Liouville integrable realisations



Another Liouville integrable realisation of Kepler

- $(H, |L|^2, L_z)$ is the standard Liouville integrable realisation of the Kepler problem
- (H, G, L_z) with $G = |L|^2 + 2ae_z$ is another Liouville integrable realisation of the Kepler problem
- a is deformation parameter, for $a=0$: $G=|L|^2$
- corresponding statements for the quantum case; use quantum Laplace-Runge-Lenz vector

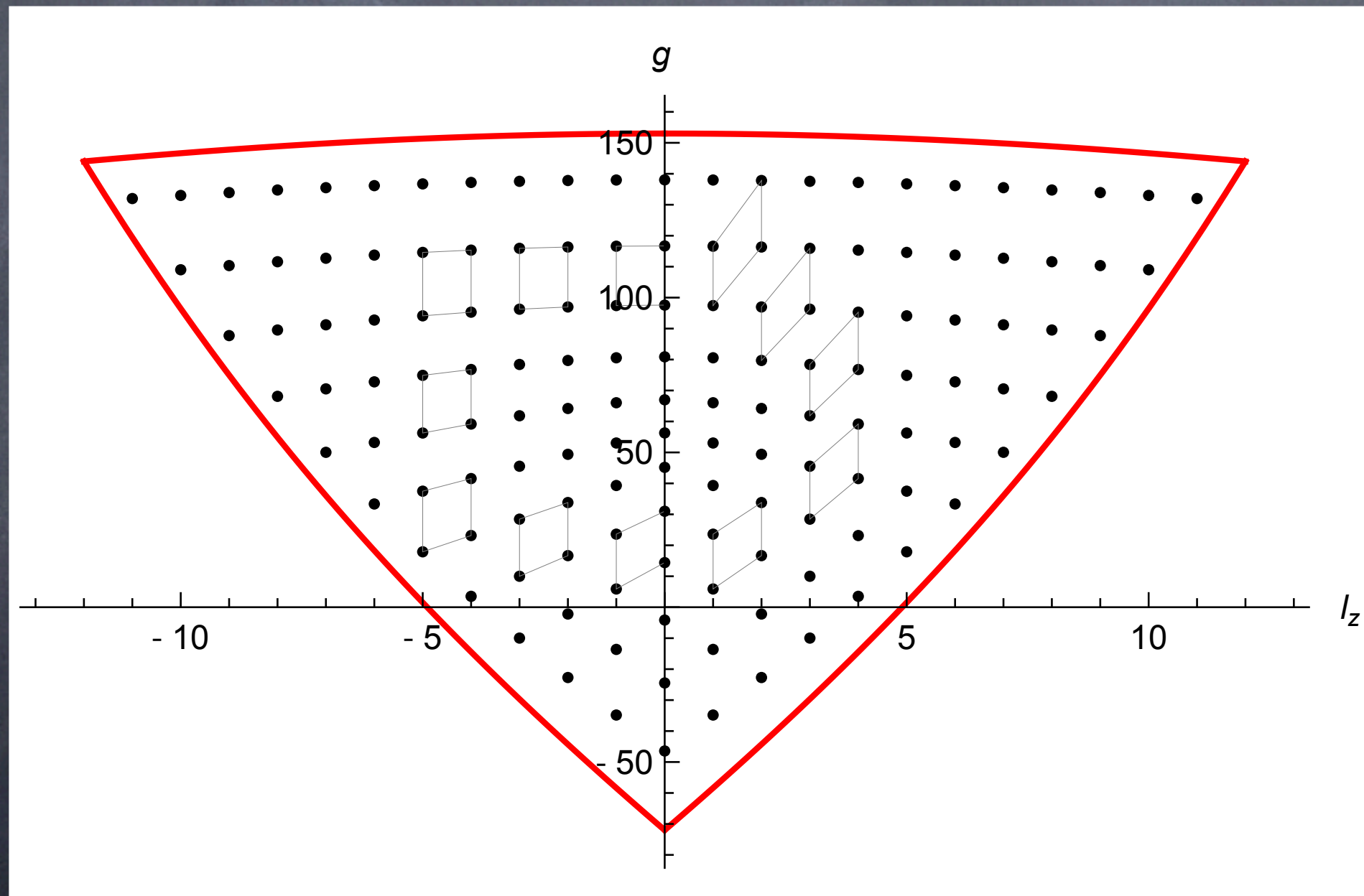


- classical: it appears arbitrary to consider Liouville integrable realisations of a super-integrable system
- quantum: each realisation should be studied since each gives a set of commuting observables, which can hence be **measured simultaneously**
- can measure $(H, |L|^2, L_z)$ simultaneously and can measure (H, G, L_z) simultaneously, but cannot measure $(H, |L|^2, G, L_z)$ simultaneously



Joint spectrum (G, L_z)

$$H = -1/(2n^2) \text{ fixed}$$



$$n=12$$

$$a=28.8$$



Monodromy in Hydrogen

- Consider the quantum integrable system (H, G, L_z) in the semiclassical limit.
- Theorem (Waalkens and HRD): For fixed n with $n^2 > a$ the joint spectrum of (G, L_z) has quantum monodromy. The joint spectrum is locally a lattice, but there is only one globally well defined quantum number m , the eigenvalue of L_z .

HRD, Waalkens, PRL 120:020507 (2018)



The Kepler system on $S^2 \times S^2$

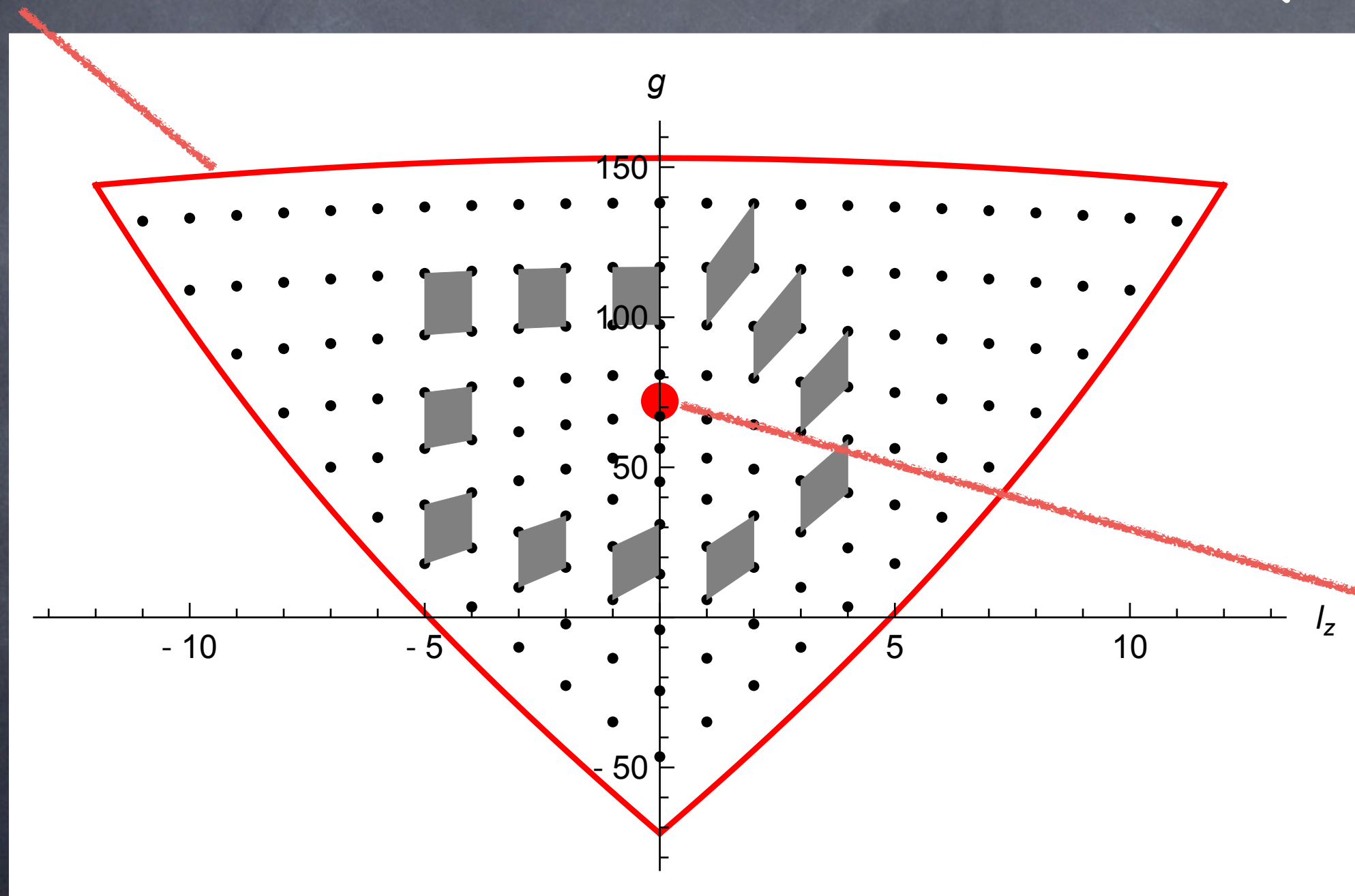
- commuting fct. in K,L-variables: $G = |L|^2 + 2aK_z/n$, L_z
- on $S^2 \times S^2$: Coordinates $S_1=(x_1,y_1,z_1)$, $S_2=(x_2,y_2,z_2)$
radius of spheres: n
 $G = |S_1+S_2|^2 + 2a(z_1-z_2)/n$, $L_z = z_1+z_2$
- commuting fct. $x_1x_2+y_1y_2+z_1z_2 + f(z_1,z_2)$ and z_1+z_2 ;
here $f(z_1,z_2) = b(z_1-z_2)$
- other systems on $S^2 \times S^2$ with different f :
 - * coupled spins: Zhilinskii & Sadovski '99, Floch & Pelayo '16, Jaume Alonso & HRD & Hohloch '18, Jaume Alonso's poster
 - * geodesic flow on S^3 : Diana Nguyen's poster
 - * Hydrogen EM field: Cushman & Sadovskii '99, & Efsthathiou '07
 - * 2FF: Hohloch & Palmer '18



Joint spectrum (G, L_z)

critical values of classical momentum map

$H = -1/(2n^2)$ fixed



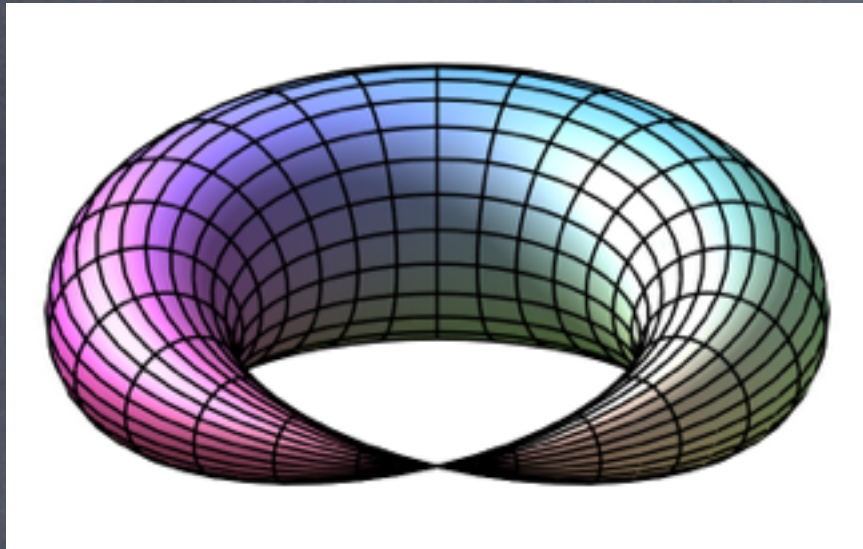
$n=12$

$a=28.8$

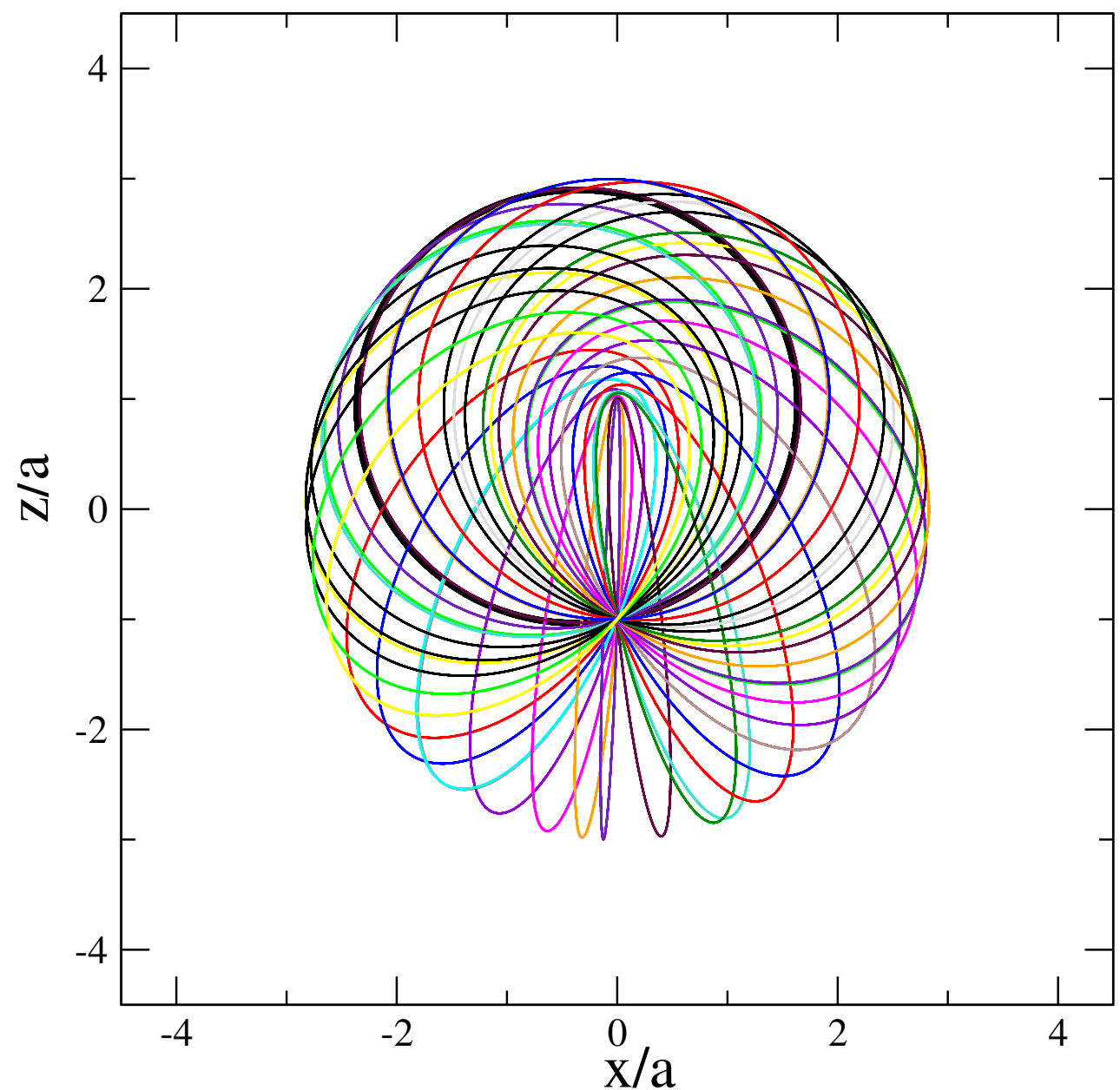
focus-
focus
value



pinched torus



pre-image of
focus-focus value

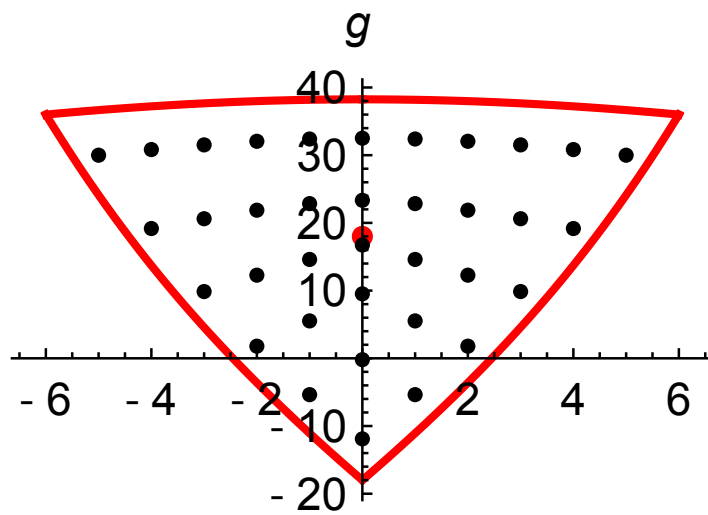


Kepler ellipses forming
the pinched torus

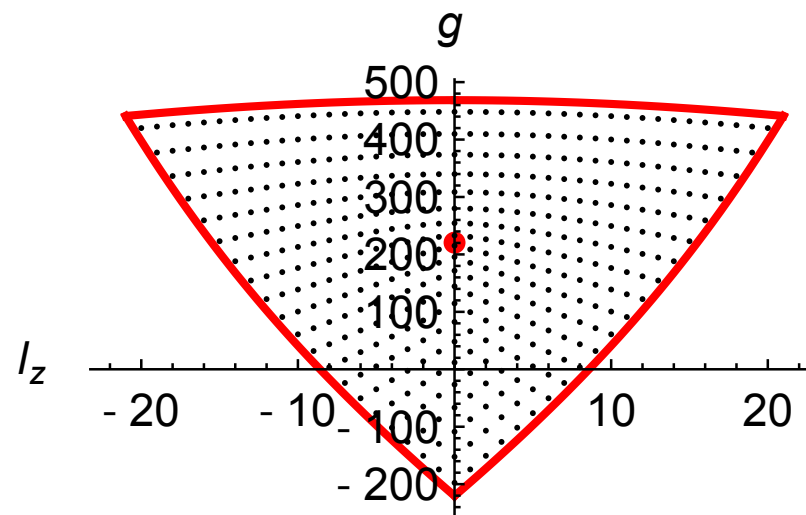
Holger Dullin



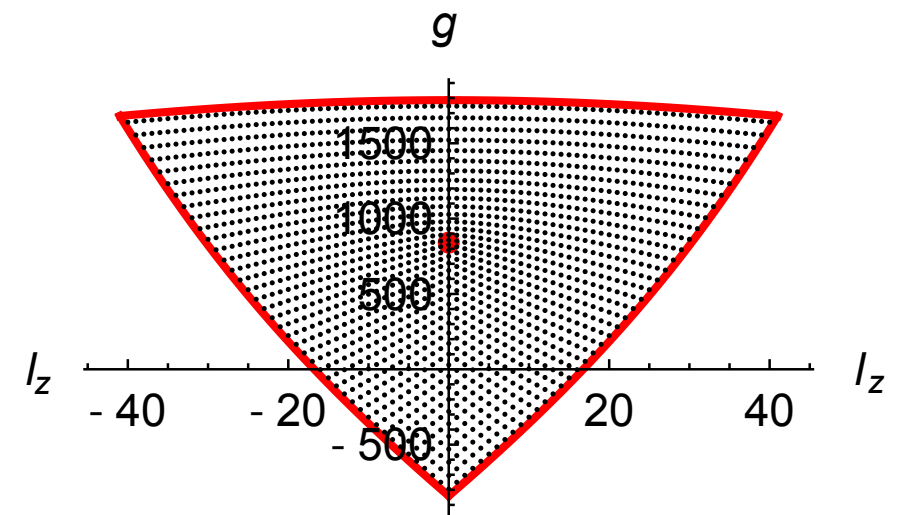
Semiclassical limit



$$n = 6$$



$$n = 21$$



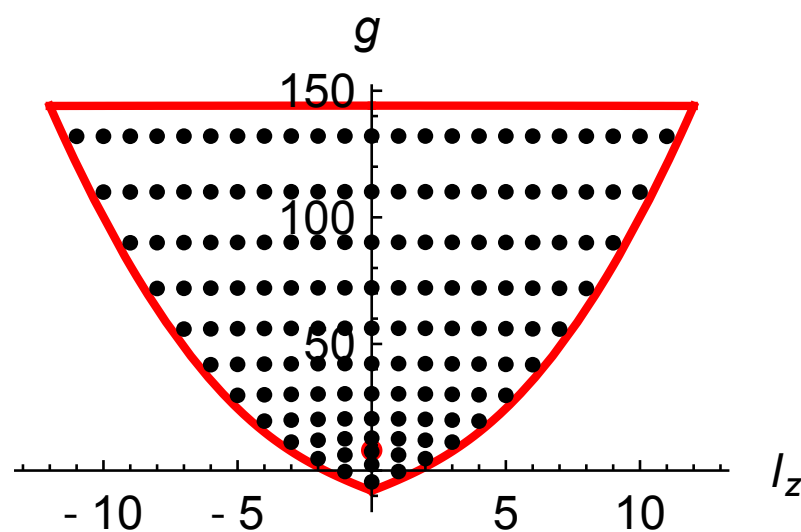
$$n = 41$$

$a=n^2/4$: keeping the relative position of focus-focus point fixed



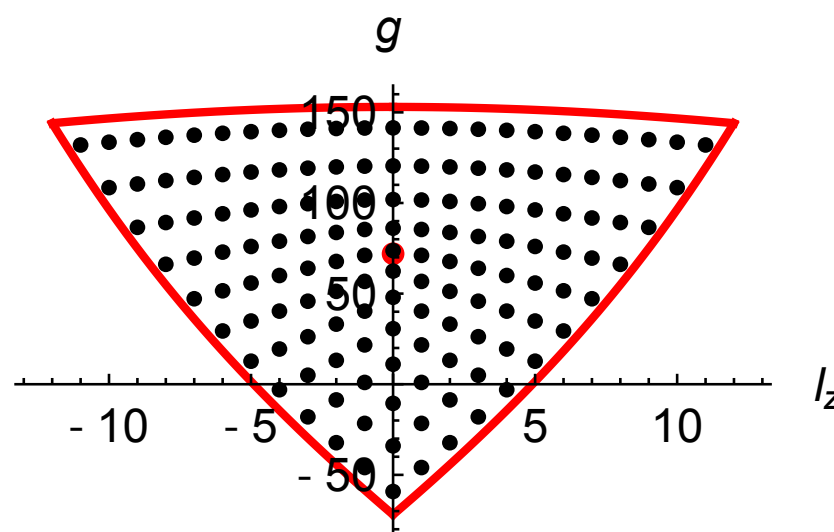
Dependence on a

$n=12$

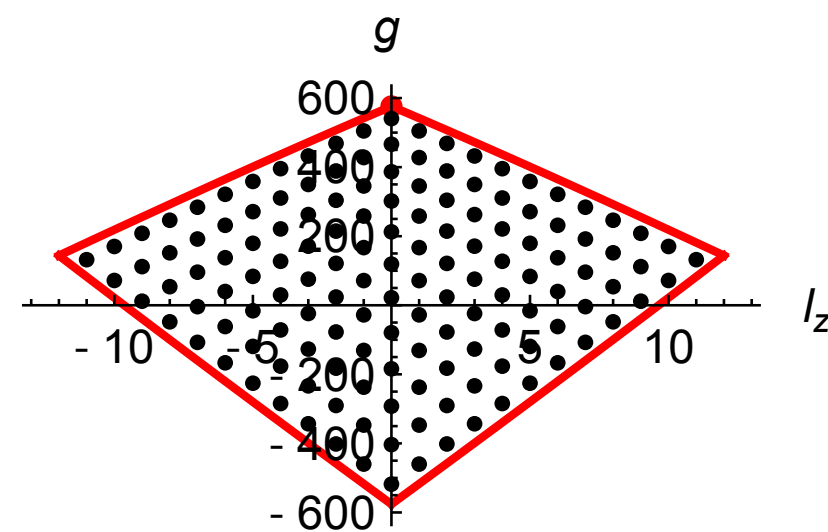


$a=4$

$n^2 > a$: monodromy,
semi-toric systems



$a=36$



$a=288$

$n^2 < a$: no monodromy,
a toric system

parabolic quantum numbers n_1, n_2, m
 $n = n_1 + n_2 + |m| + 1$

Holger Dullin



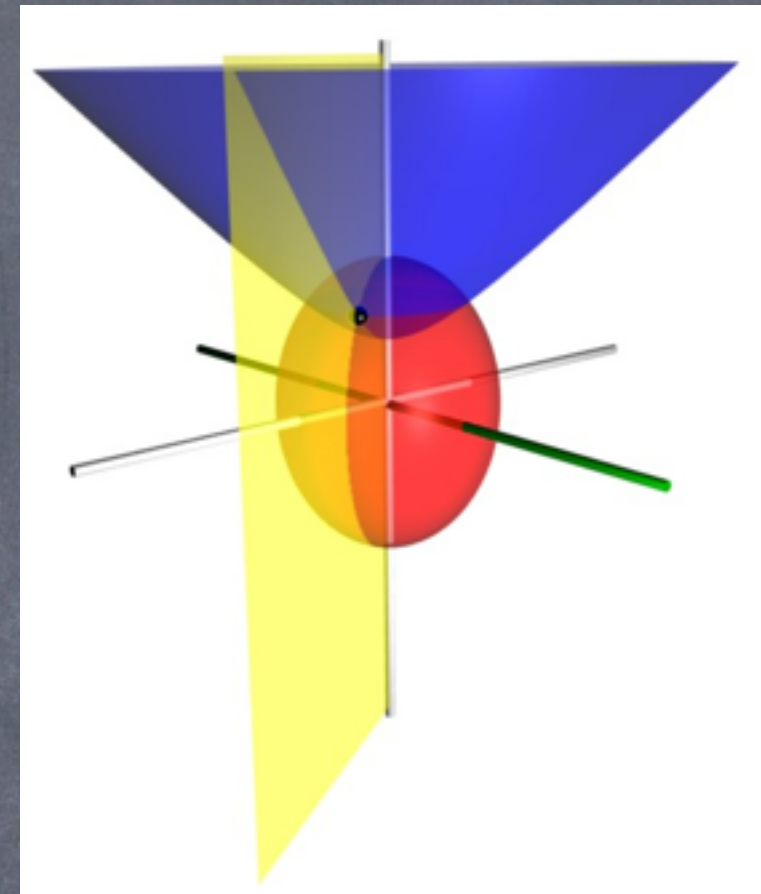
multi-separability

- separability in different coordinate systems; implies super-integrability
- gives natural realisations of Liouville integrable systems, similarly in quantum mechanics
- Kepler problem (Hydrogen) is multi-separable
 - 1) spherical coordinates ($a=0$)
 - 2) parabolic coordinates ($a=\infty$)
 - 3) spheroidal coordinates
 - 4) sphero-conical coordinates
- The family (H, G, L_z) contains the first three



prolate spheroidal coordinates

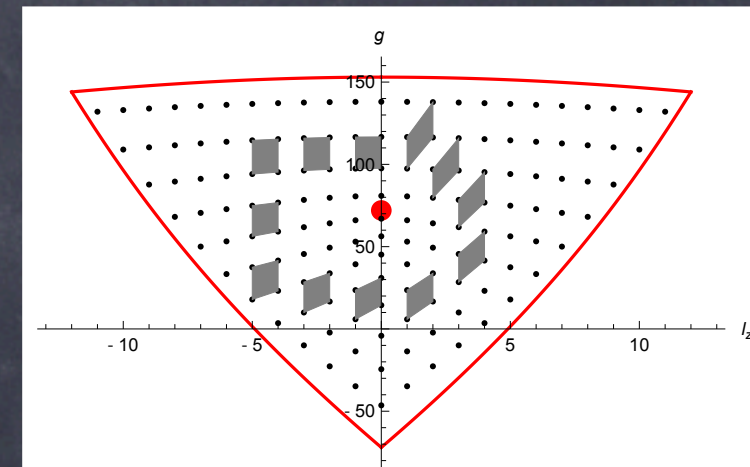
- confocal ellipses, focus points at $\pm a$
- classical: separation constant G , as before $G = |L|^2 + 2ae_z$
- quantum: separation gives spheroidal wave equation (confluent Heun)
- finite expansion of spheroidal harmonics in terms of spherical harmonics (Coulson & Robinson '58)

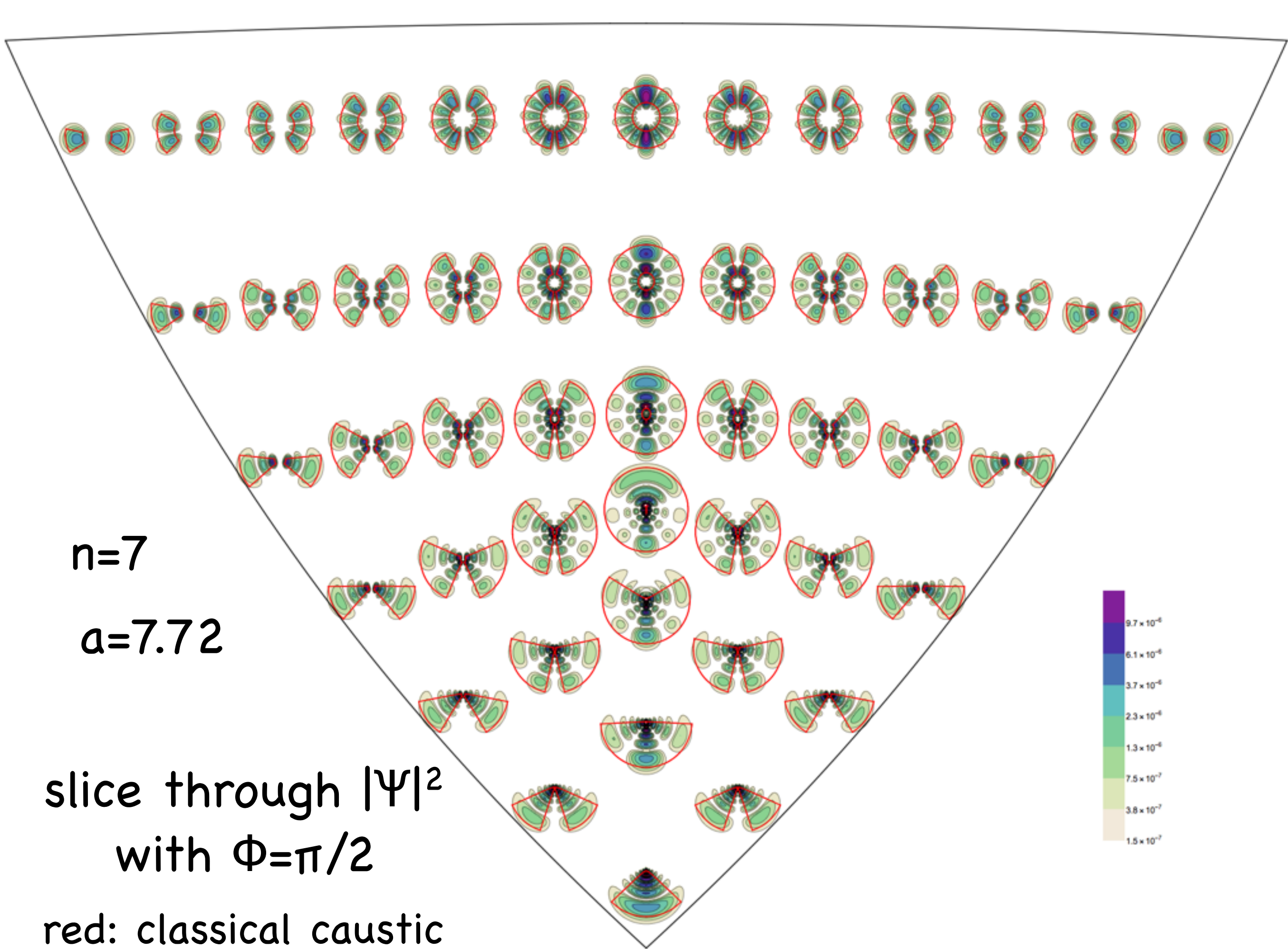


Spectrum and Orbitals

- Classical and Quantum system on $S^2 \times S^2$
- Obtained from quotient by flow of H
- Joint spectrum determined by eigenvalues of tri-diagonal matrix A , fixed eigenvalue m of L_z
- non-zero matrix entries for $l = |m|, |m|+1, \dots, n-1$ for fixed n, m

$$A_{l,l} = l(l+1), \quad A_{l+1,l} = \frac{a}{n} \sqrt{(n^2 - l^2)(l^2 - m^2) / (l^2 - \frac{1}{4})}$$





Regularisation

- Moser regularisation: maps system to geodesic flow on S^3 . Clearly $SO(4)$ symmetric. Quotient by flow of H leads to system on $S^2 \times S^2$.
- Kustaanheimo–Stiefel regularisation: maps system to isotropic harmonic oscillator in \mathbb{R}^4 . Quotient by flow of H and flow of bi-linear integral leads to system on $S^2 \times S^2$.
- Liouville integrable realisations obtained from separation of variables are the same as before.



Laplacian

- remove the potential \rightarrow just free particle,
QM = Laplace operator
- separates in many coordinate system
- separated equations define special functions
spherical harmonics, spheroidal harmonics,...
- quotient by free particle dynamics
leads to integrable systems on T^*S^2
- with Diana Nguyen and Sean Dawson:
Monodromy in spheroidal harmonics, Sean's poster



free particle quotient

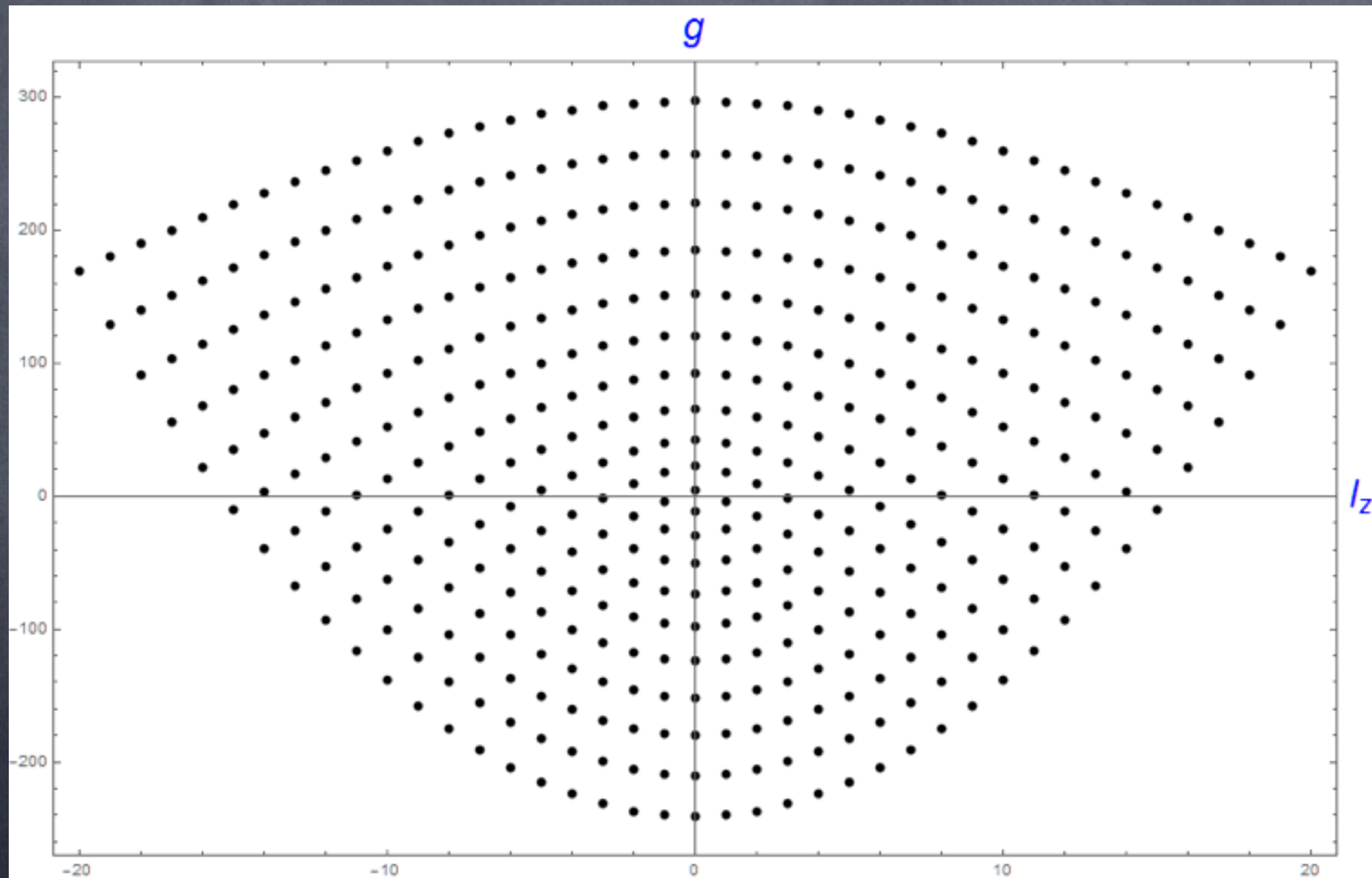
- $r = (x, y, z)^{\dagger}$, $p = (p_x, p_y, p_z)^{\dagger}$, $L = r \times p$ ang. momentum
- $H = (p_x^2 + p_y^2 + p_z^2)/2$, flow of H : straight lines in \mathbb{R}^3
- Fix H : sphere in momentum space
- Components of linear momentum p and angular momentum $L = r \times p$ satisfy $SE(3)$ algebra
- reduced phase space is T^*S^2
- spheroidal integrable system:
 $G = |L|^2 - a^2(p_x^2 + p_y^2)$, L_z



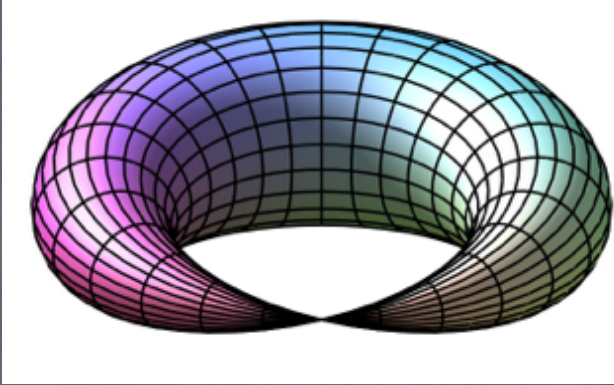
Monodromy in Spheroidal Harmonics

with Diana Nguyen and Sean Dawson

```
ListPlot[ Flatten[ Table[{l, SpheroidalEigenvalue[n, l, 16]}, {n, 0, 20}, {l, -n, n}], 1]]
```



Proof of Monodromy



- Regular reduction by flow of H to $S^2 \times S^2$
- Singular reduction by flow of L_z to 2-D orbifold using invariant polynomials
- Show existence of a pinched torus by reconstruction from orbifold to $S^2 \times S^2$
- similar to Cushman & Sadvskii '00, & Efsthathiou '04, but our "Hamiltonian" G is different
- implies classical monodromy (Matveev '96, Zung '97)
- implies quantum monodromy (Vu Ngoc '99)



singular reduction

- reduction by flow generated by L_z , rotation

- invariant polynomials:

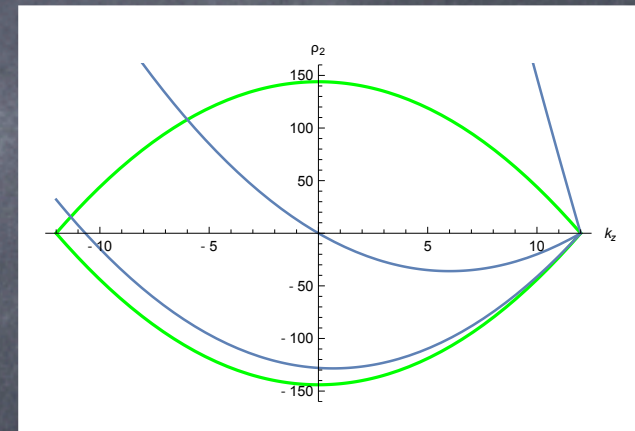
$$I_1 = K_z, \quad I_2 = L_x^2 + L_y^2 - K_x^2 - K_y^2, \quad I_3 = K_x L_y - K_y L_x,$$

- $\{I_1, I_2 + iI_3\} = 2i(I_2 + iI_3), \quad \{I_2, I_3\} = \partial C / \partial I_1$

$$\text{Casimir } C = P(I_1) - I_2^2 - I_3^2,$$

- $G = I_2 + n^2 + m^2 - I_1^2 + 2aI_1/n$

- $C=0$ defines reduced phase space, singular for $m = L_z = 0$ (Lemon), reconstruction gives pinched torus



Conclusion

- Super-integrable systems have different Liouville-integrable realisations, some may have monodromy
- Same for corresponding Quantum Integrable Systems
- The most important super-integrable systems:
 - * Kepler problem (PRL 120:020507 (2018))
 - * 3D isotropic Harmonic Oscillator
 - * Free particle (Laplace Operator) (Sean's poster)
 - * geodesic flow on sphere (Diana's poster)
- All have (quantum) monodromy in spheroidal coordinates
- Program: Identify semi-toric special functions

