

A scenic mountain landscape at sunrise or sunset, with a hiker's silhouette in the foreground. The sun is low on the horizon, casting a warm glow over the rugged, rocky terrain. The hiker is positioned on the right side of the frame, leaning forward and holding a climbing tool. The background shows a range of mountains under a clear sky.

Using enhanced sampling methods to  
calculate free energy barriers

A generic approach to kinetics?

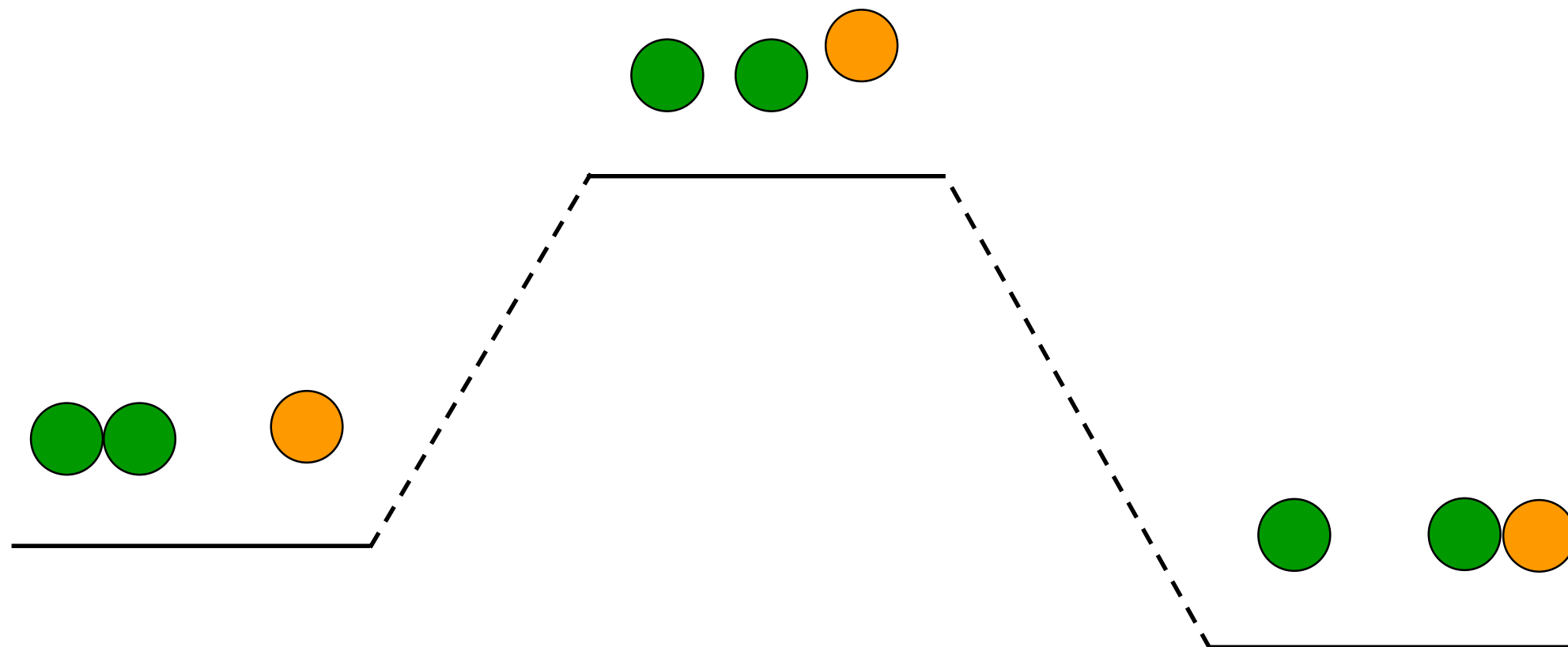
Dr. Kristof Bal

Virtual seminar @ Vrije Universiteit Brussel, 15 October 2020

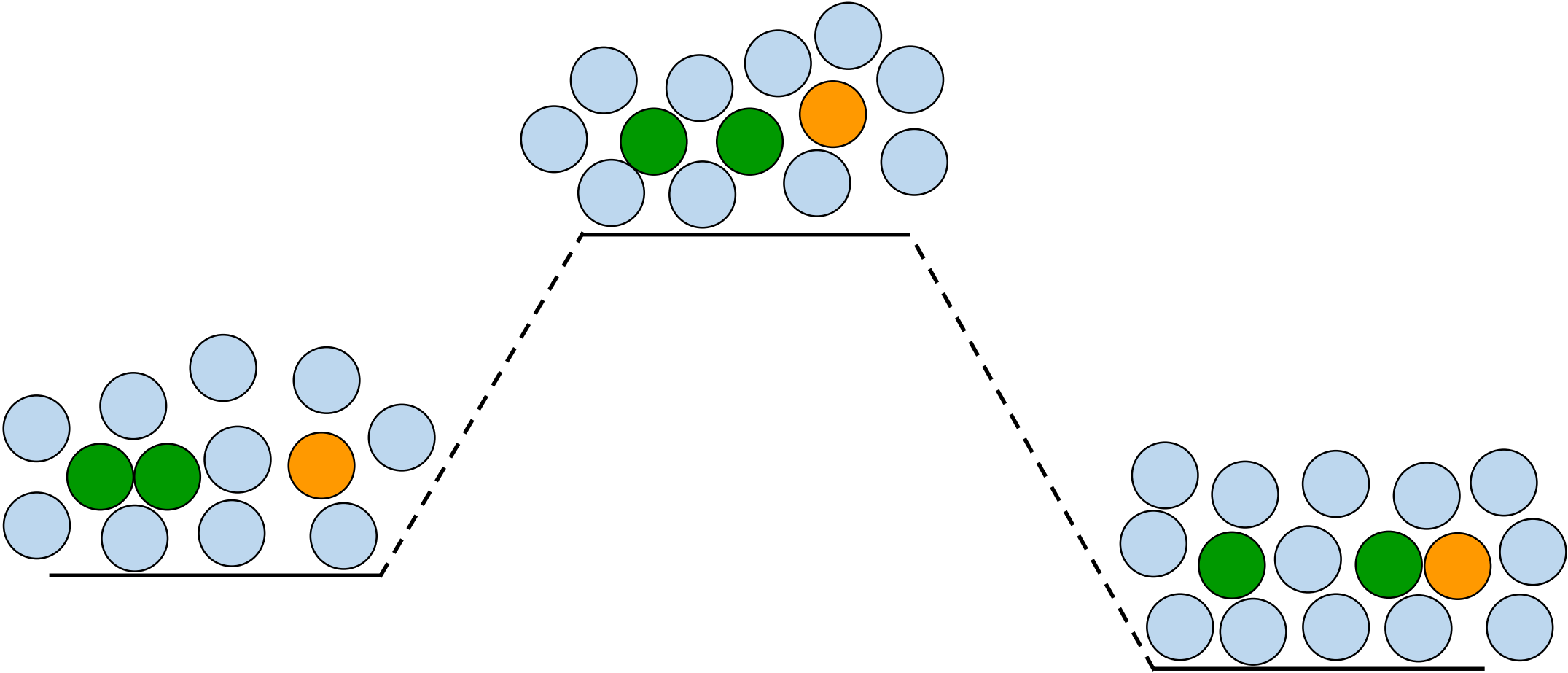
**Setting the scene**

Free energy, collective variables, enhanced sampling

In many cases, a “static” model of chemistry suffices



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Or not!

Most chemical problems are actually about ensemble properties.

→ **Molecular dynamics simulations** can sample many configurations

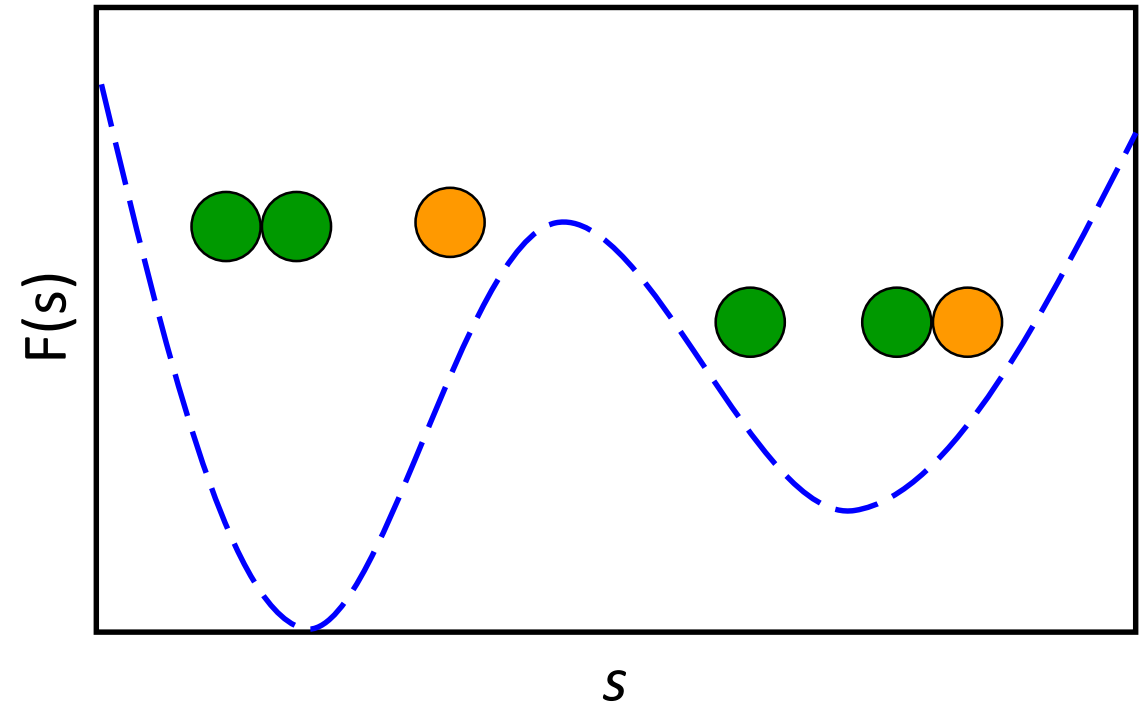
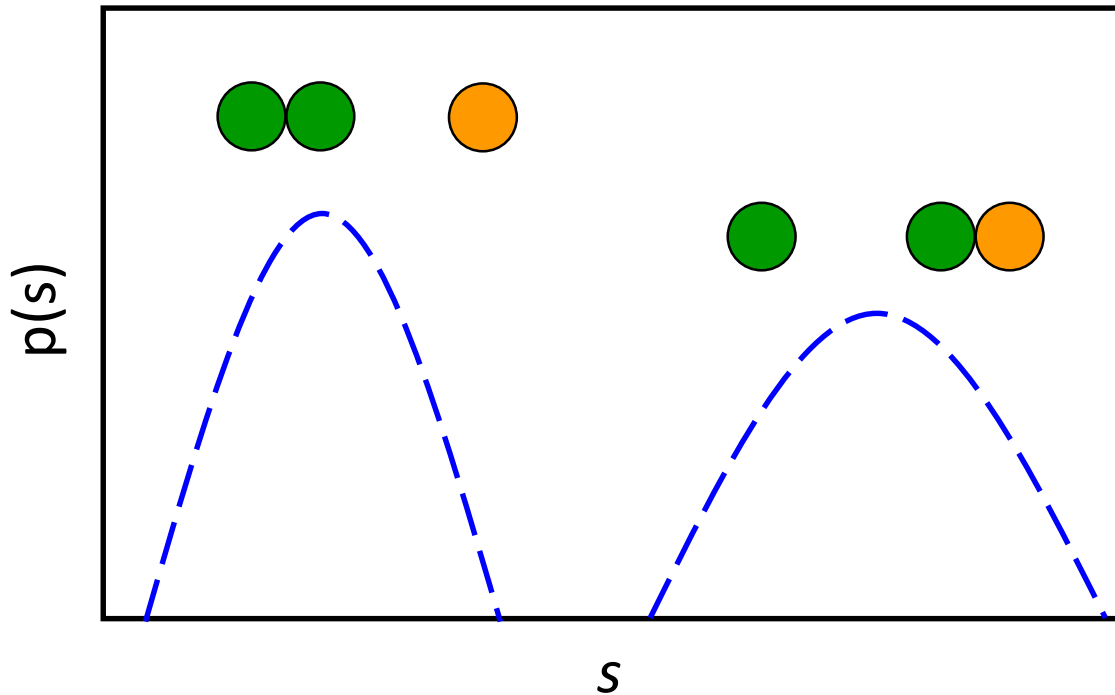
MD simulations produce high-dimensional data – hard to analyse!

→ Reduce dimensionality of the problem using **collective variables**

Set of CVs  $\mathbf{s}$ : order parameters that can distinguish between relevant states of the system (preferably intuitively)

$$p(\mathbf{s}) = \langle \delta[\mathbf{s} - \mathbf{s}(t)] \rangle$$

$$F(\mathbf{s}) = -\frac{1}{\beta} \ln p(\mathbf{s})$$



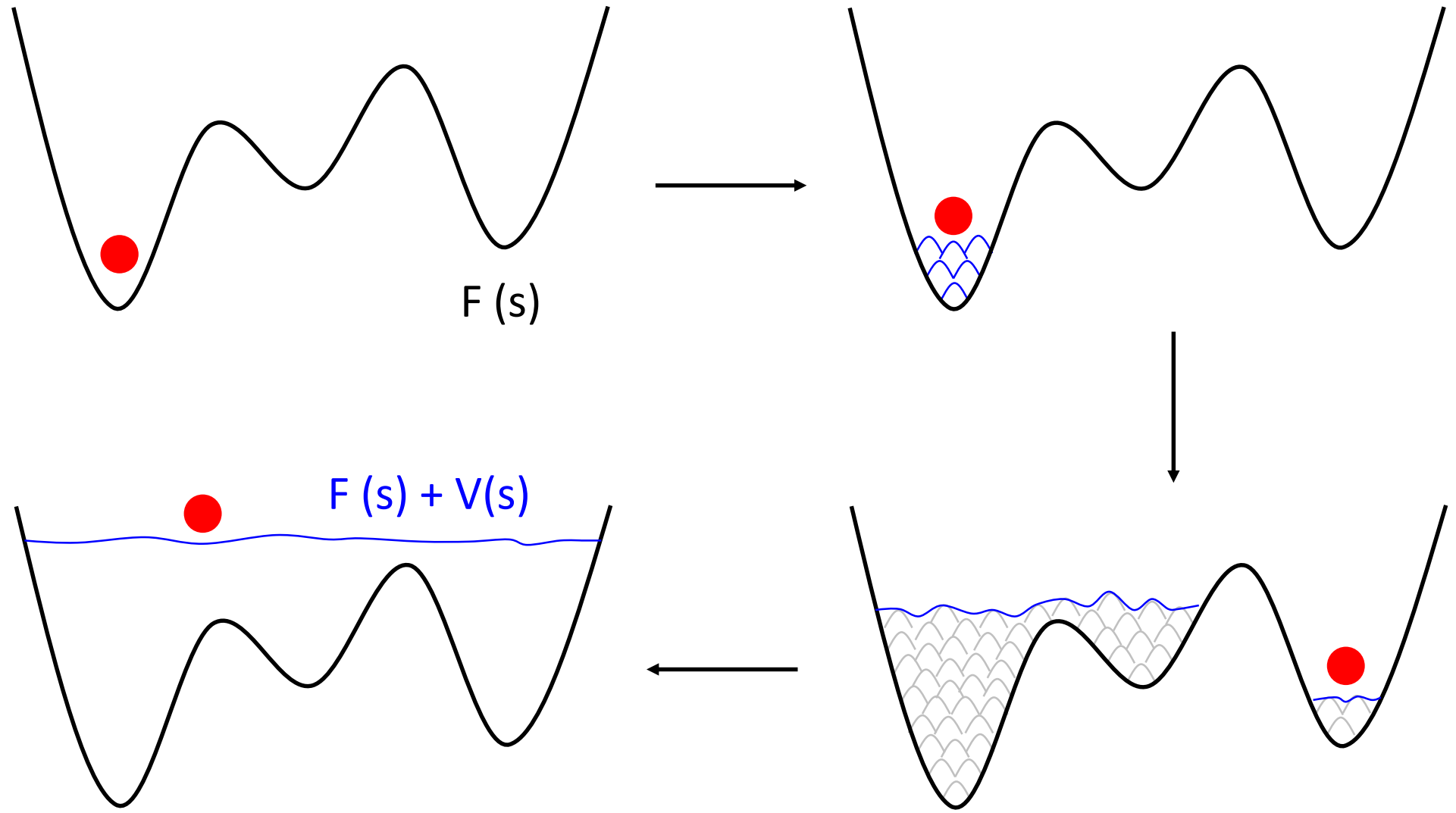
The free energy surface  $F(\mathbf{s})$  in principle encodes all thermodynamic information of a transformation along  $\mathbf{s}$ .

But in practice it's hard to efficiently sample the histogram  $p(\mathbf{s})!$

→ System remains trapped in metastable states during (short) MD time scale

Need for **enhanced sampling methods**, e.g., metadynamics

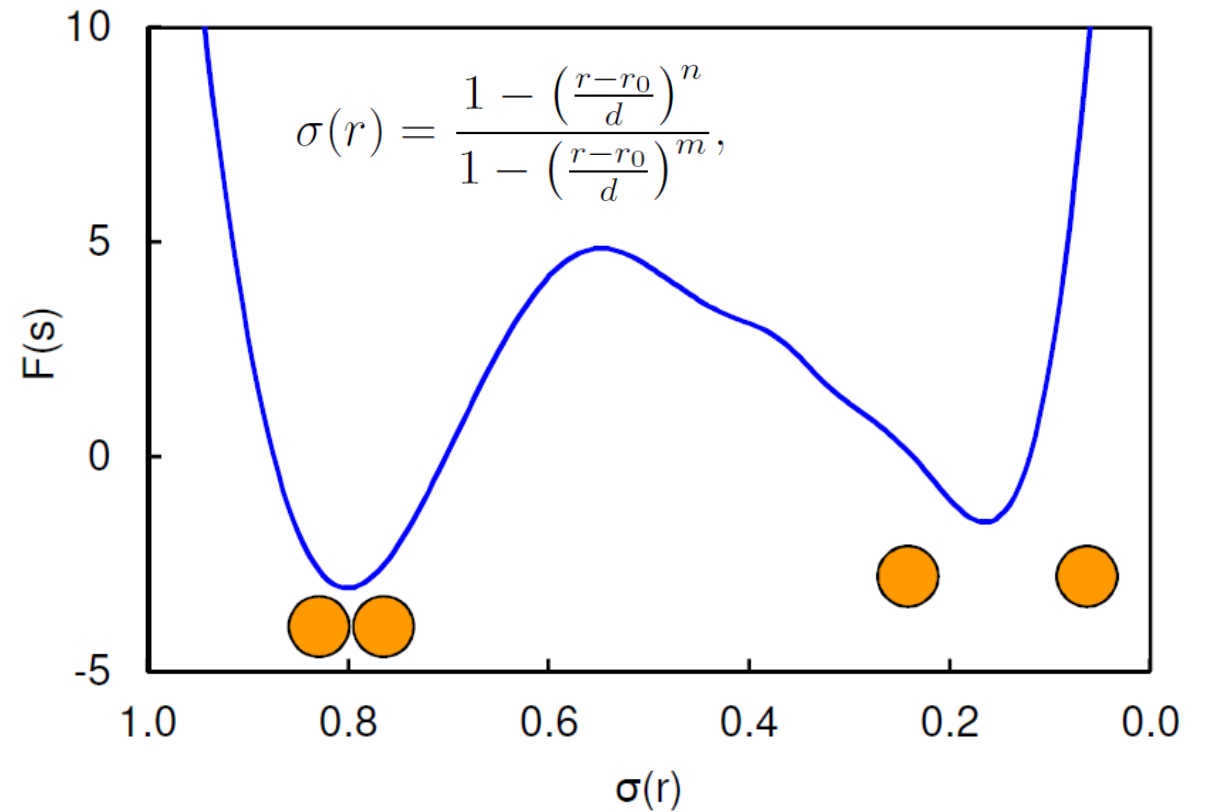
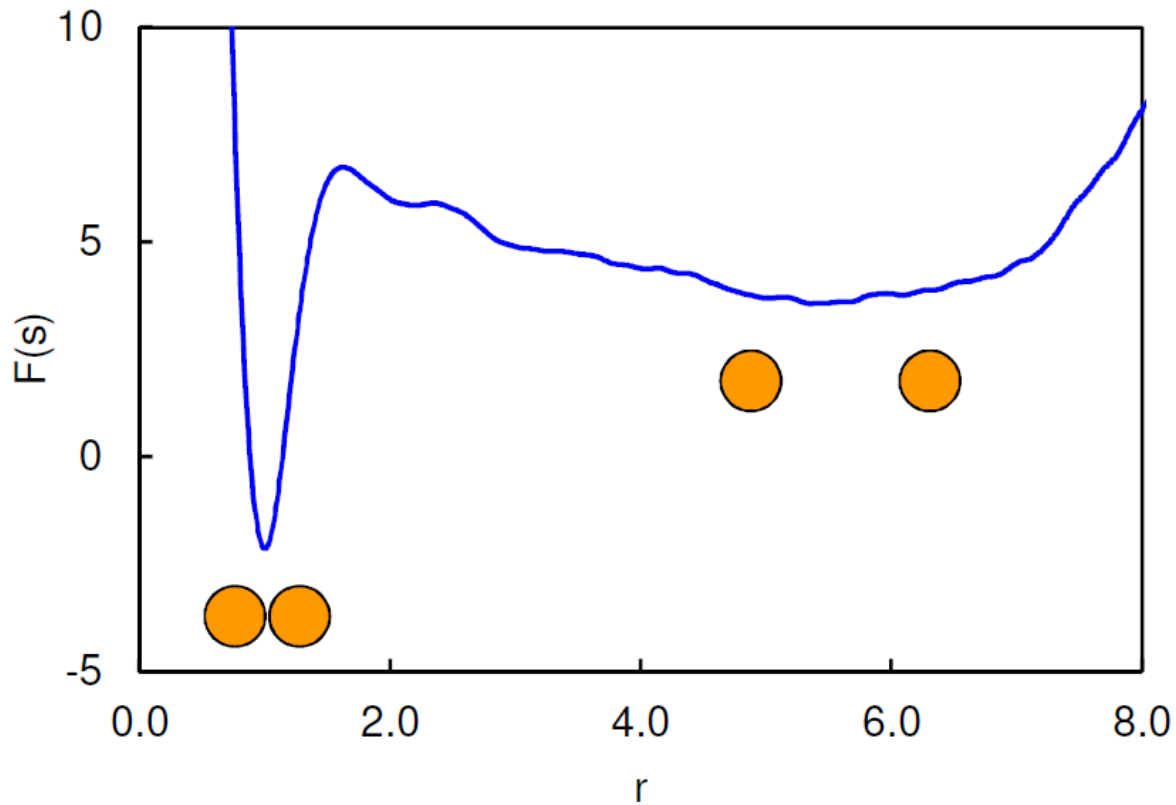
# Metadynamics: add bias potential $V(s)$ to accelerate sampling





Need to reweight the histogram:  $p(\mathbf{s}) = \langle w \cdot \delta[\mathbf{s} - \mathbf{s}(t)] \rangle_b$

Even for a single degree of freedom, we can choose many parametrizations of  $s$ .



Not an issue: integrate the probability density  $p(\mathbf{s})$  (per unit of  $\mathbf{s}$ )

$$F_A = -\frac{1}{\beta} \ln \int_A d\mathbf{s} e^{-\beta F(\mathbf{s})},$$

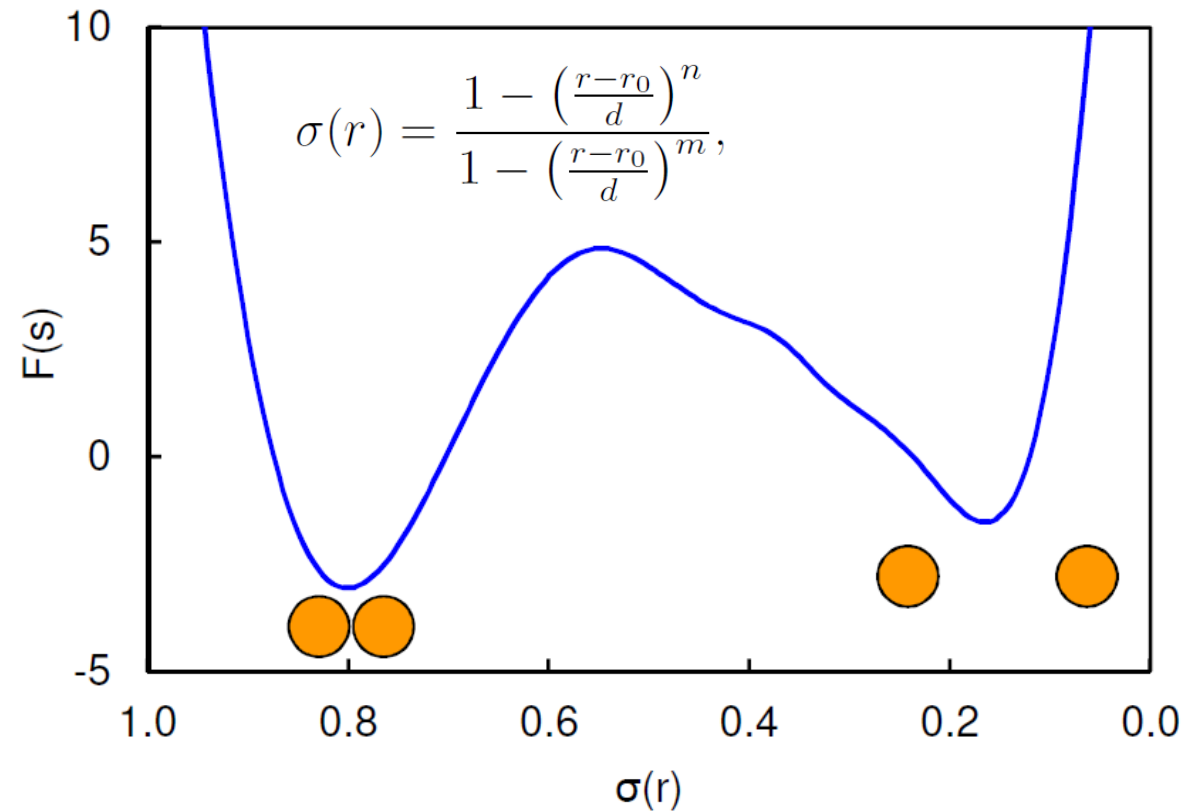
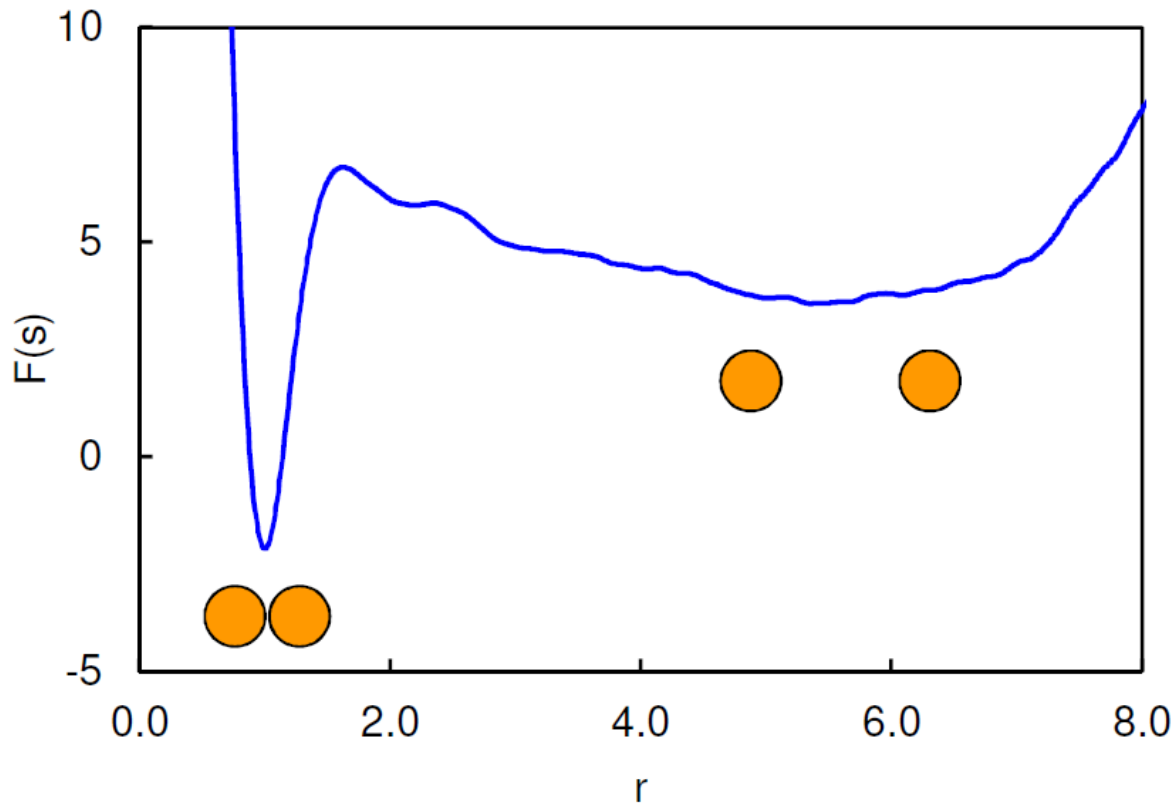
$$F_B = -\frac{1}{\beta} \ln \int_B d\mathbf{s} e^{-\beta F(\mathbf{s})},$$

$$\Delta F_{A \rightarrow B} = F_B - F_A.$$

This way, free energy differences between states (and equilibrium constants) can be consistently computed

For this system (dimer in Lennard-Jones fluid), both CVs do give a consistent **binding free energy**  $\Delta F = -1.4$

So... how about the **free energy barrier**  $\Delta^\ddagger F$ ?



## **Kinetics from the free energy - Pitfall 1**

What is the free energy of the transition state?

$F_{\text{TS}} = F(\mathbf{s}_{\text{TS}})$  is not the free energy of the transition state, because  $p_{\text{TS}} = p(\mathbf{s}_{\text{TS}})$  is a probability **density** which is not invariant to  $\mathbf{s}$ .

So  $F_{\text{TS}} - F_A$  is a meaningless quantity, and not the barrier!

It turns out that:

$$\Delta^\ddagger F_{A \rightarrow B} = F_{\text{TS}} + \frac{1}{\beta} \ln \frac{\langle |\nabla s| \rangle_{s_{\text{TS}}}^{-1}}{h} \sqrt{\frac{2\pi m}{\beta}} - F_A$$

$$\Delta^\ddagger F_{A \rightarrow B} = F_{\text{TS}} + \frac{1}{\beta} \ln \left( \frac{\langle |\nabla s| \rangle_{s_{\text{TS}}}^{-1}}{h} \sqrt{\frac{2\pi m}{\beta}} \right) - F_A$$

We see the **partition function in the segment  $s$**

Essentially, we have to take out the degree of freedom along  $s$

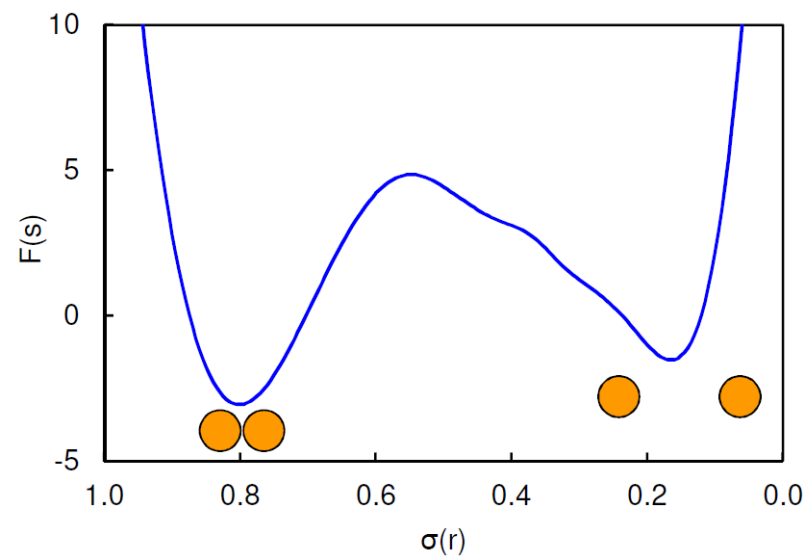
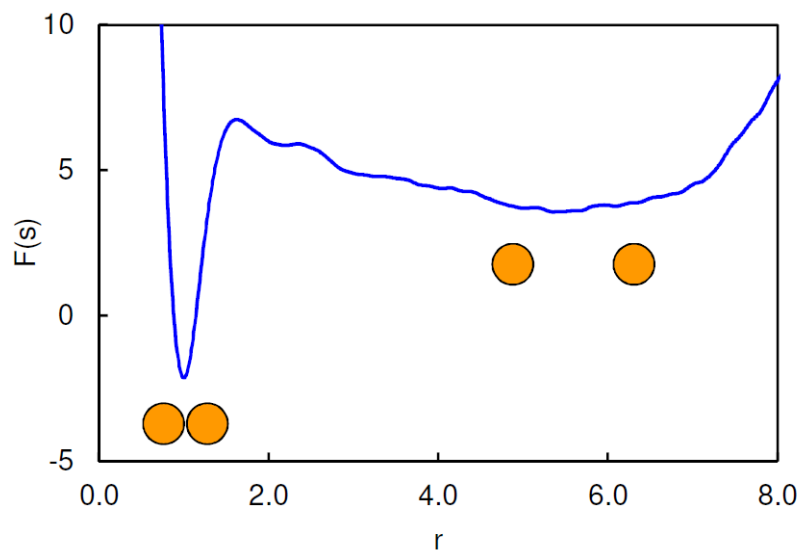
because there is no density along this direction

We can use transition state theory (TST) to get rates

$$k_{A \rightarrow B} = (h\beta)^{-1} e^{-\beta \Delta^\ddagger F_{A \rightarrow B}}$$

# Consistency!

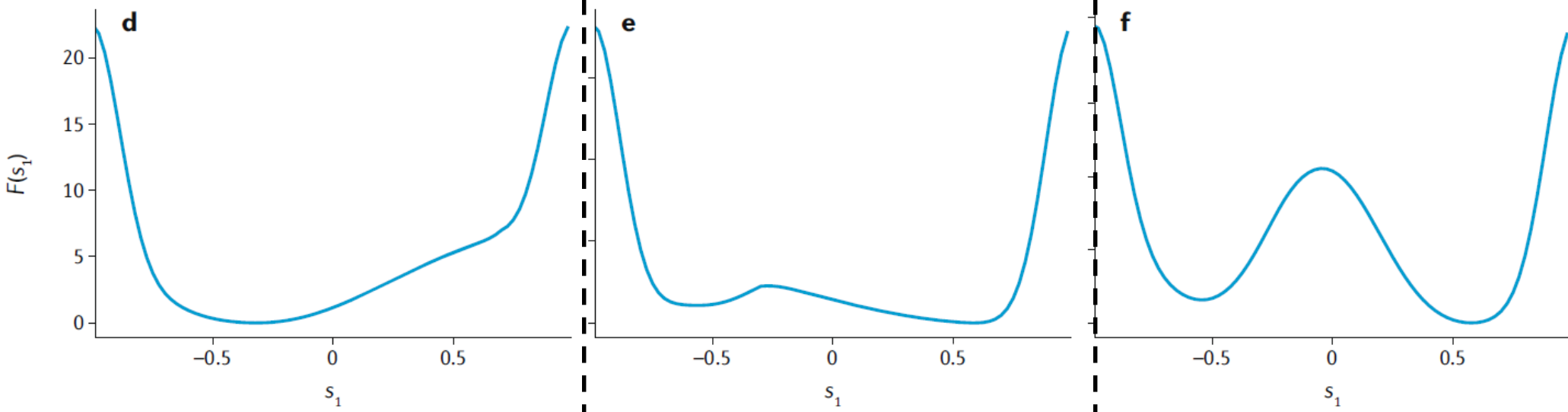
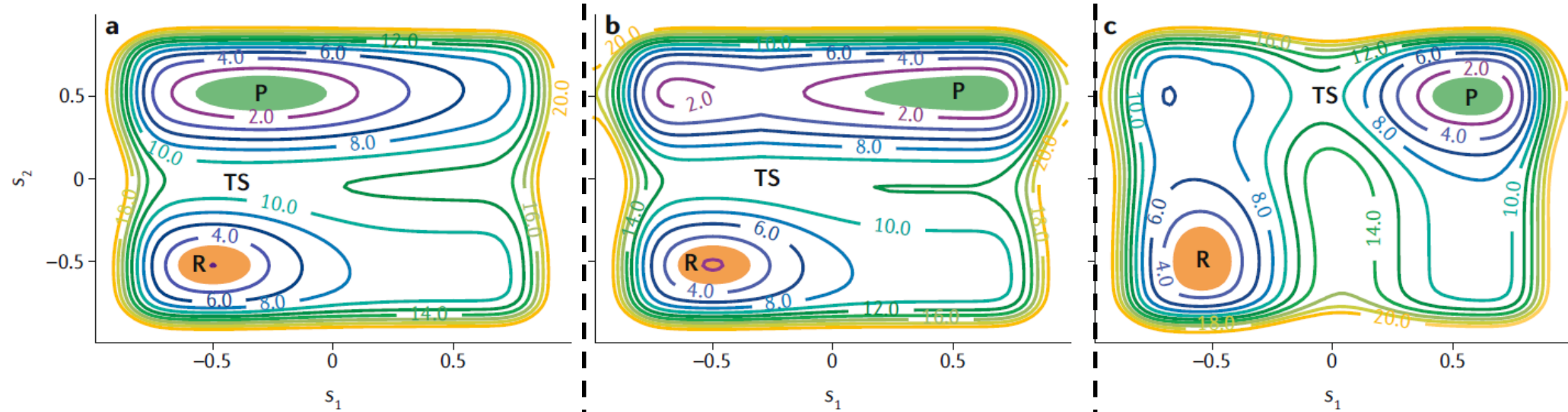
energy/CV	r	$\sigma(r)$
$\Delta F$	-1.39	-1.38
$F_{\text{TS}} - F_{\text{A}}$	5.04	3.15
$\Delta^\ddagger F$	3.44	3.41



## **Kinetics from the free energy - Pitfall 2**

Is our CV actually the reaction coordinate?





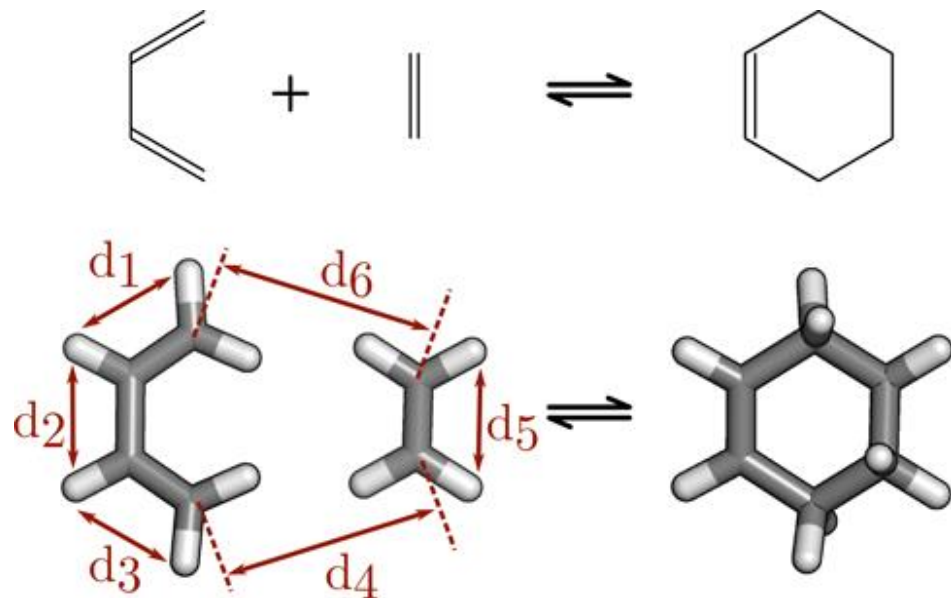
useless!

only thermodynamics  
inefficient biasing

good barrier  
best for biasing

Many new methods promise to produce good CVs systematically:

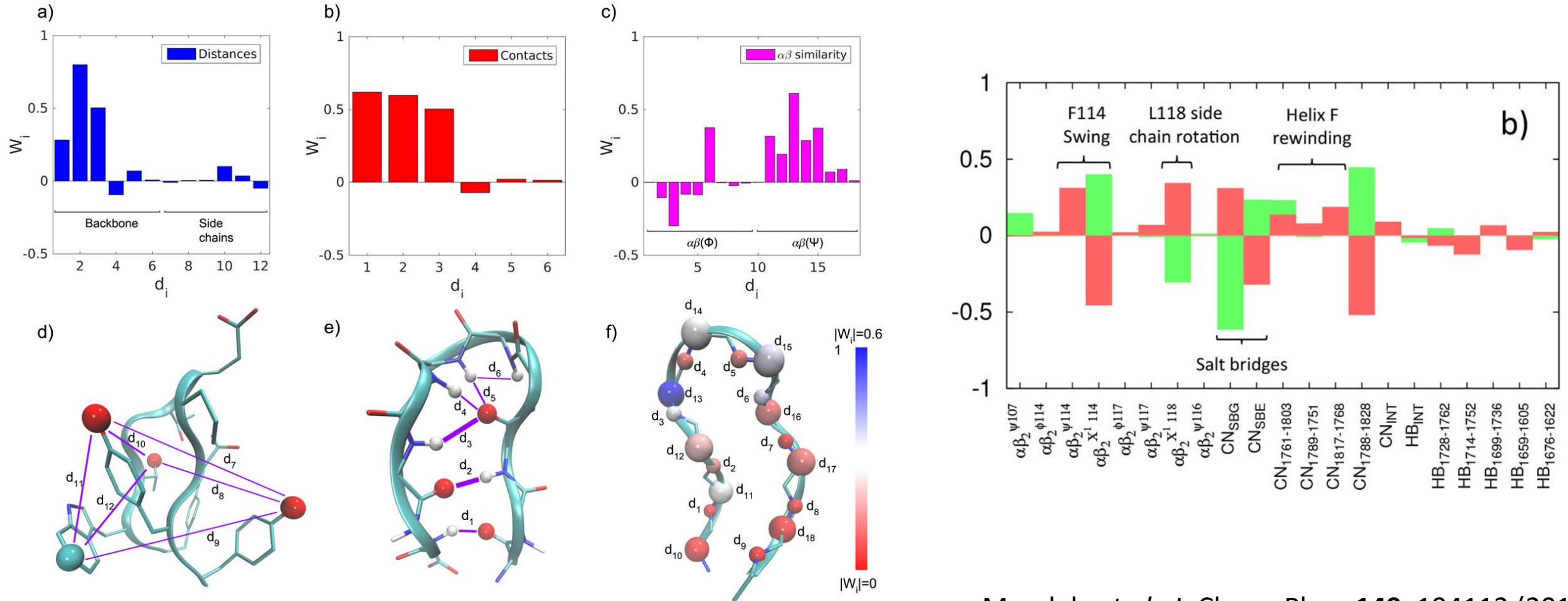
1. Start with large set of candidate CVs  $\mathbf{s} = (s_1, s_2, s_3, \dots)$
2. Assume true RC =  $f(\mathbf{s})$  (linear expansion, neural network, ...)
3. Construct *a priori*, or refine iteratively during simulation



$d_1$	$d_2$	$d_3$	$d_4$	$d_5$	$d_6$
0.05	-0.21	0.08	-0.69	0.00	-0.68

# Need for good CVs: bug or feature?

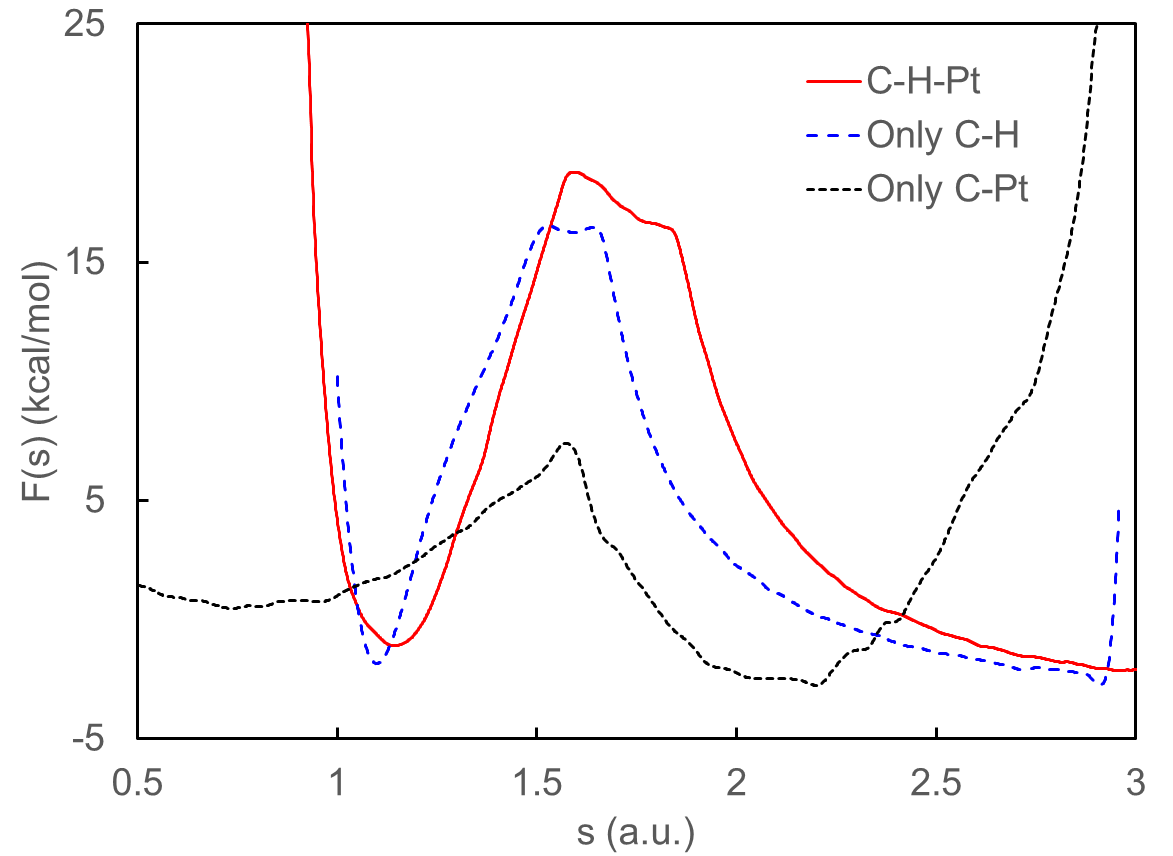
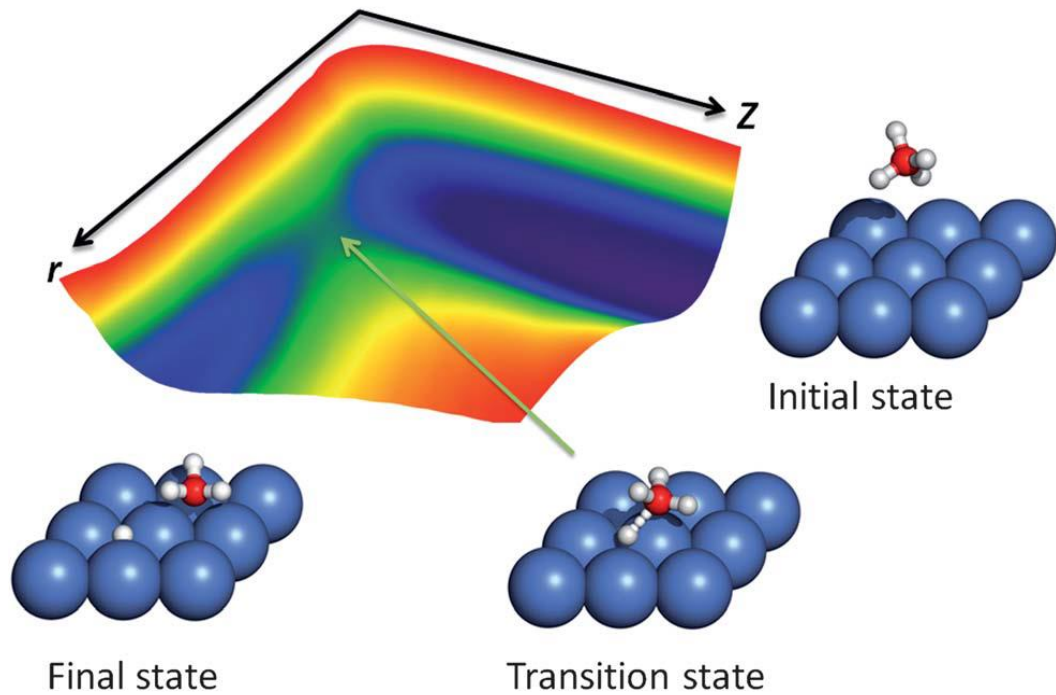
Good CV = good sampling, but also physical insight!



Mendels *et al.*, J. Chem. Phys. **149**, 194113 (2018)

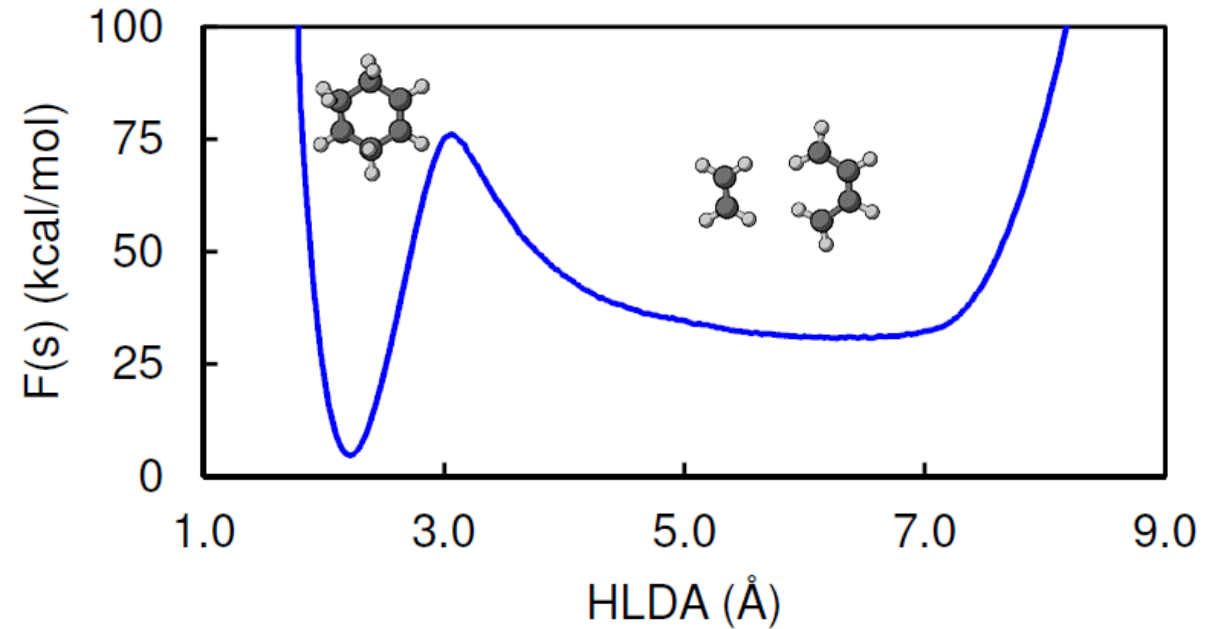
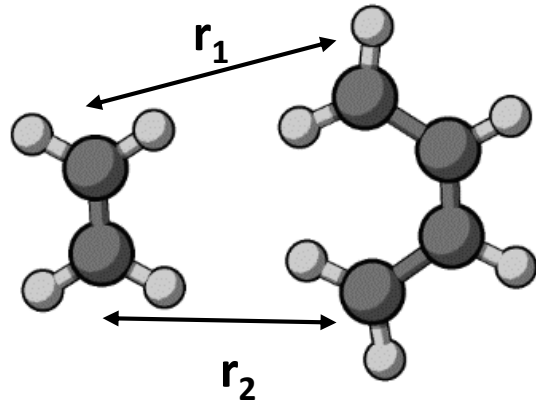
Brotzakis & Parrinello, J. Chem. Theory Comput. **15**, 1393 (2019)

# Constructing a CV for CH<sub>4</sub> chemisorption @ Pt(111)



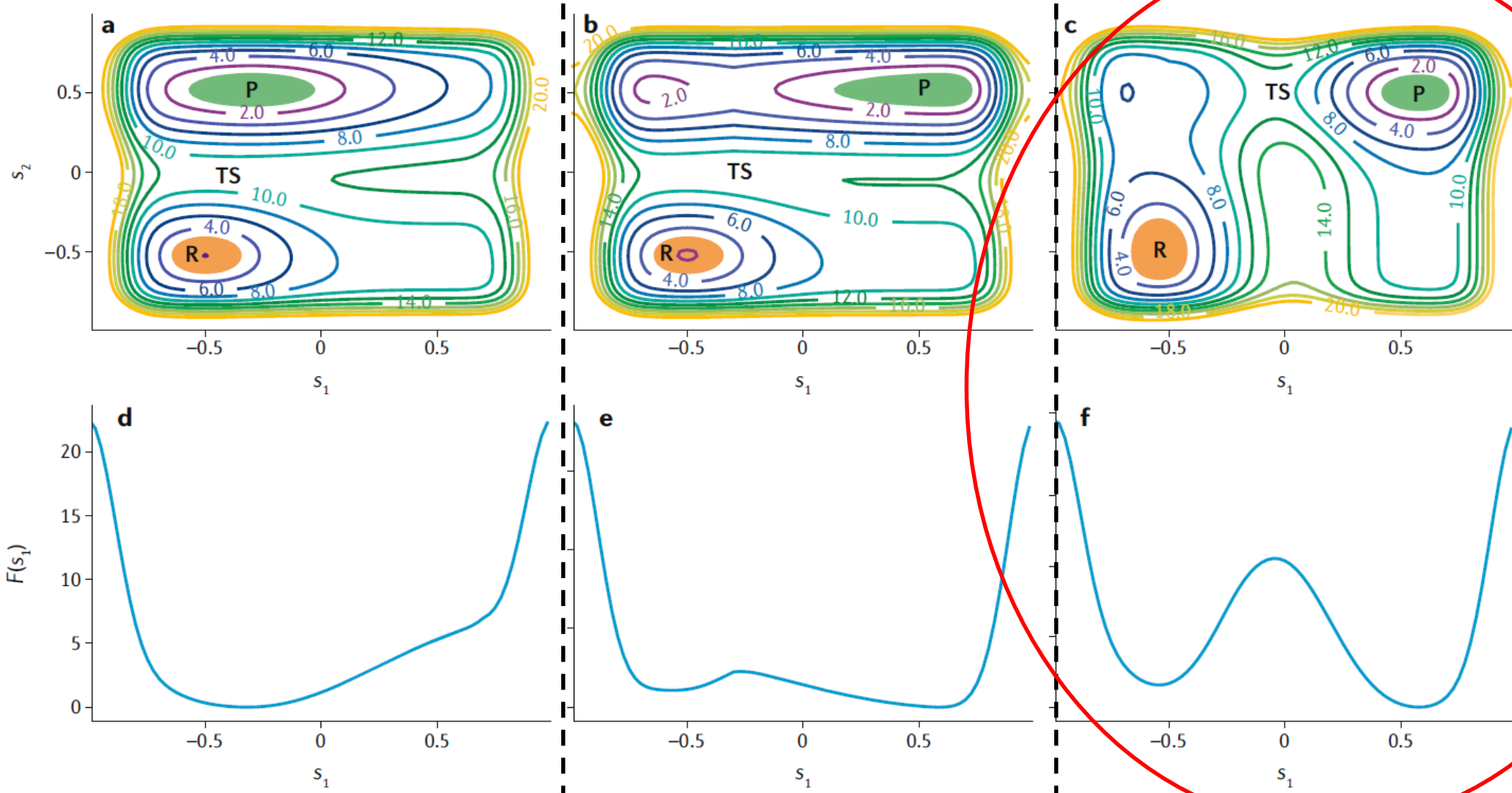
TST give the classical upper bound of the true rate

So: the best CV maximizes the barrier



Project FES on several trial CVs

- Two body vs. manybody
- 1D vs. 2D



useless!

only thermodynamics

good barrier

inefficient biasing

best for biasing

	HLDA	$r_1$	SPRINT1	SPRINT2	$\bar{\sigma}$	$(r_1, r_2)$
$\Delta F$	-23.25	-23.34	-23.25	-23.25	-23.25	-23.25
$\Delta^\ddagger F$	47.01	29.01	32.66	46.51	47.16	47.23

Two body  $r_1$  is an all-round bad CV

Several manybody CVs are decent 1D reaction coordinate

If we “give up” and go 2D, we get the best (highest) barrier

Ultimately, things are quite robust w/r/t the CV

## How good are rates?

- We (indirect):  $(7.8 \pm 1.7) \cdot 10^{-5} \text{ s}^{-1}$
- Direct approach (Fu *et al.*):  $(5.8 \pm 0.7) \cdot 10^{-5} \text{ s}^{-1}$

Direct approach = explicit sampling of reactions (with bias)

## But we need way shorter MD runs!

- We: 5 ns (250 ps for quick 'n' dirty)
- Fu *et al.*: up to 2  $\mu\text{s}$

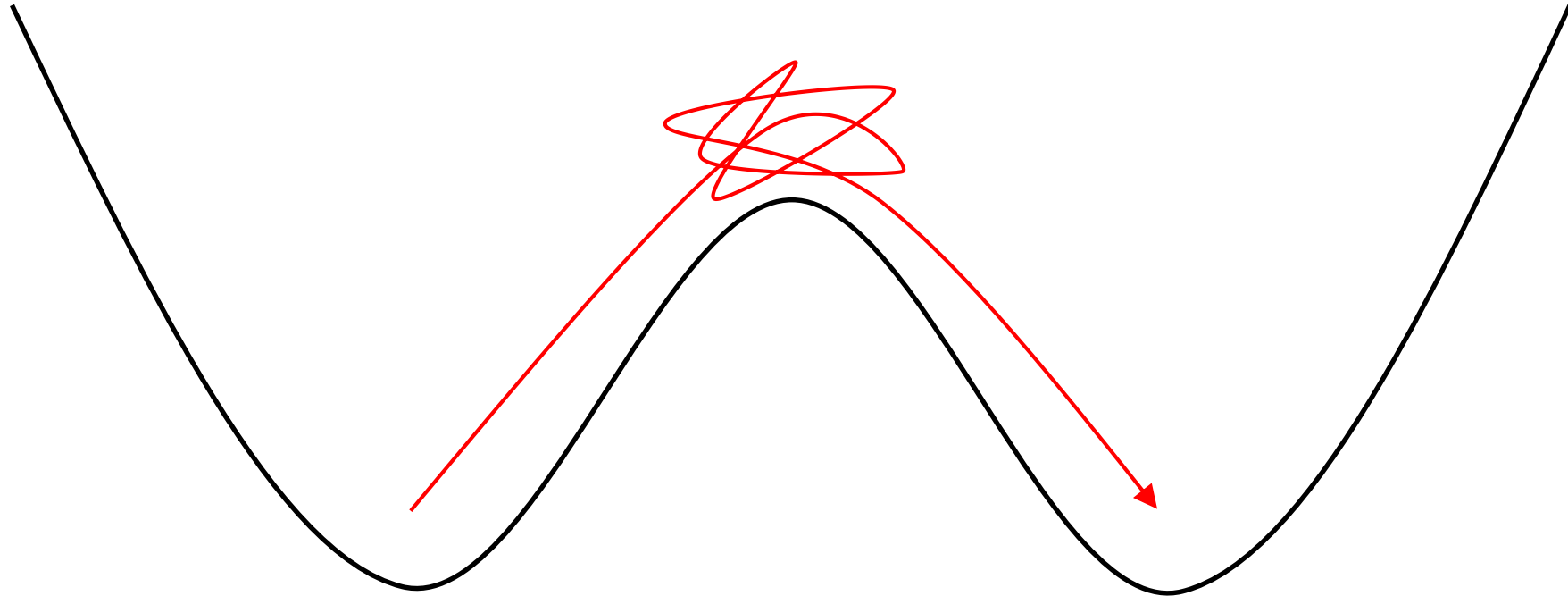
In principle we could go *ab initio*!



## **Kinetics from the free energy - Pitfall 3**

Do we move in a straight line?

TST rate = frequency of crossing TS in a certain direction



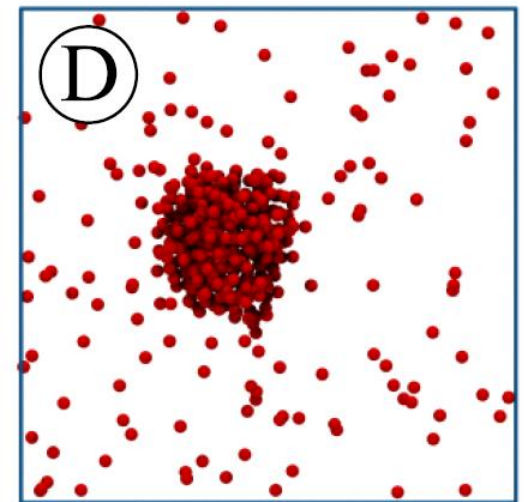
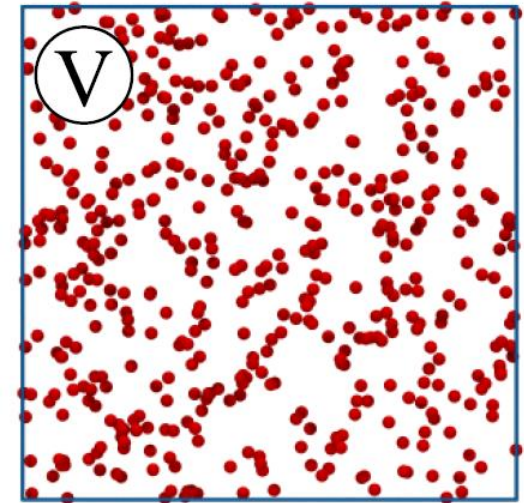
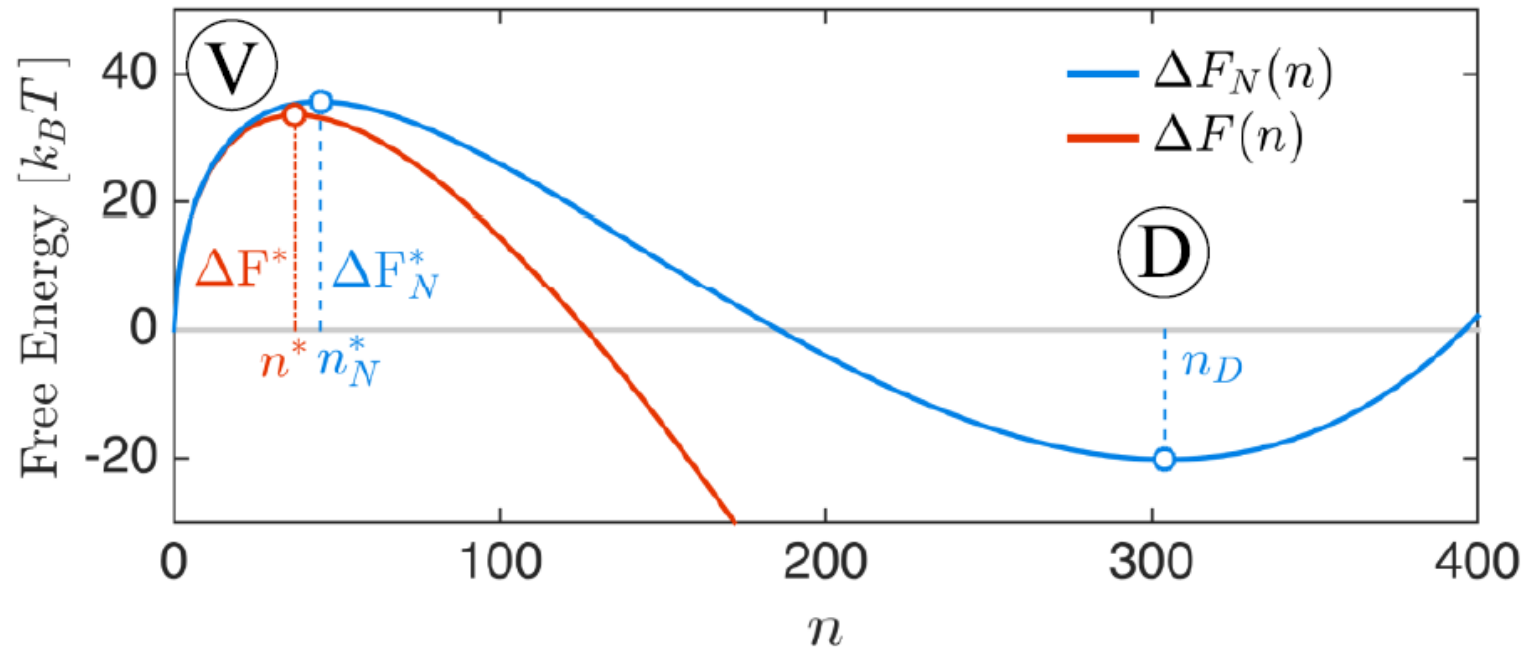
Not every TS crossing corresponds to an actual transition.

Correct the Eyring equation with a transmission coefficient  $\kappa \leq 1$

$$k_{A \rightarrow B} = \frac{\kappa}{h\beta} e^{-\beta \Delta^\ddagger F_{A \rightarrow B}}$$

# Droplet (D) nucleation from supersaturated Ar vapor (V)

No way that TST holds, right?

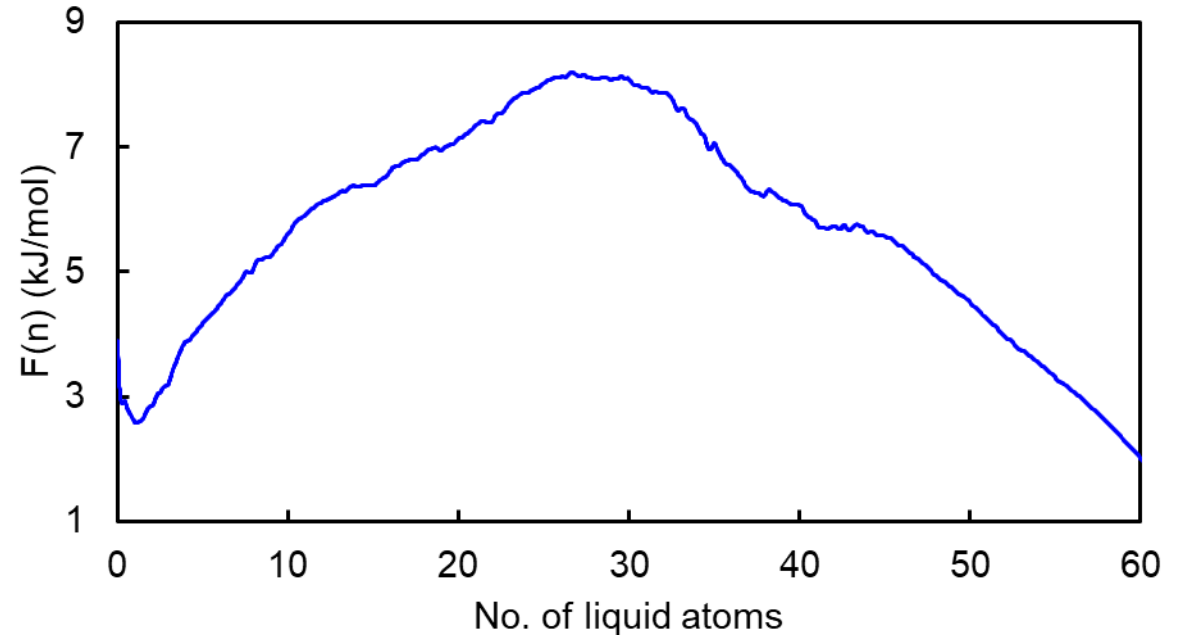


$\Delta^\ddagger F = 6.3$  kJ/mol for nucleation @ 80.07 K (supersaturation of 9.87)

### Average time to nucleation

- Direct approach: 690 ns
- TST: 8 ns

So  $\kappa \approx \mathbf{0.012}$



Short MD runs starting at TS show about 70 crossing attempts

before a liquid state is finally reached:  $\kappa = \mathbf{0.014 \pm 0.004}$

It works... and orders of magnitude faster than direct approach!

**Wrapping up**

Conclusions

## Conclusions

Studying kinetics using free energy approaches is almost as generic as studying thermodynamics (and can be done simultaneously)

Some challenges, but:

1. We learn more about our process of interest
2. Mostly shared with the field at large (so active development)

Applications only limited by imagination, including **yours**

# Acknowledgments

Thanks to Satoru Fukuhara for starting this

I took simulation ideas and set-ups from Christopher Fu, Pablo Piaggi, GiovanniMaria Piccini, Dan Mendels, Matteo Salvalaglio, ...

Erik Neyts for the support

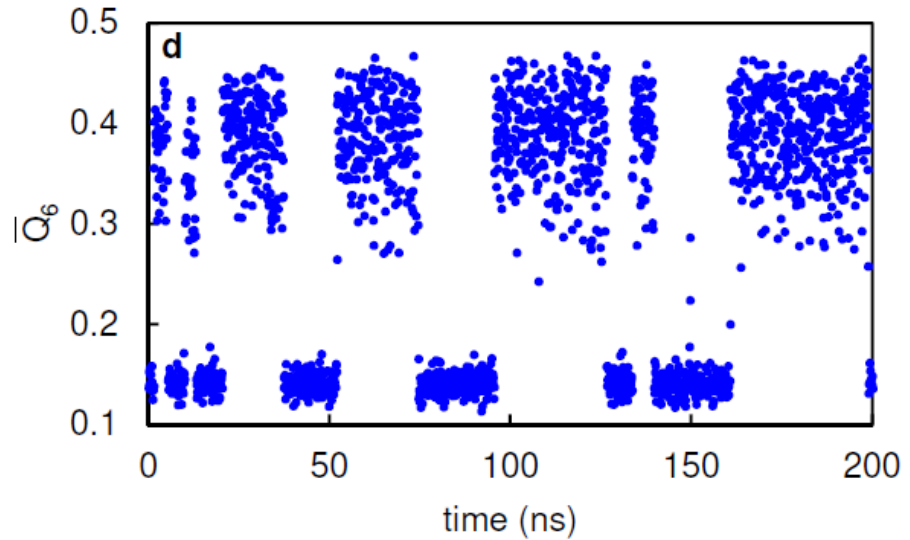
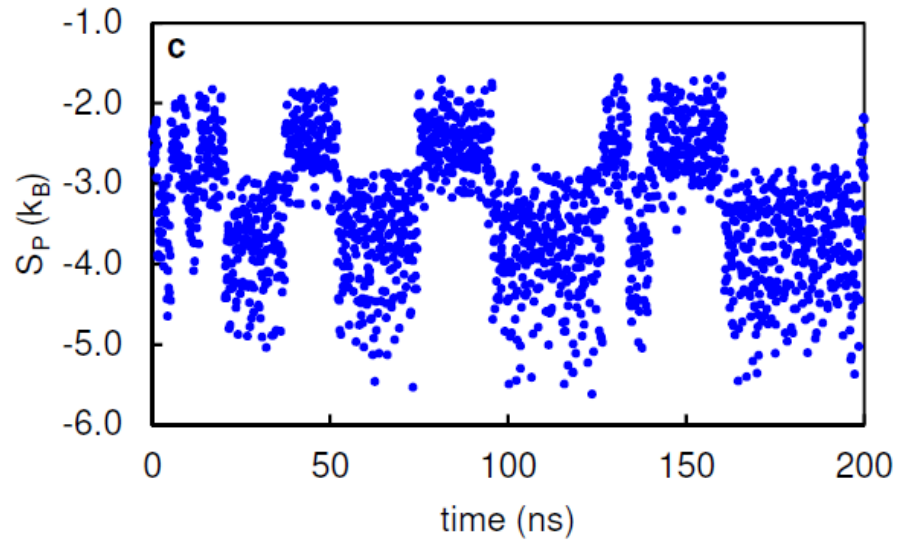
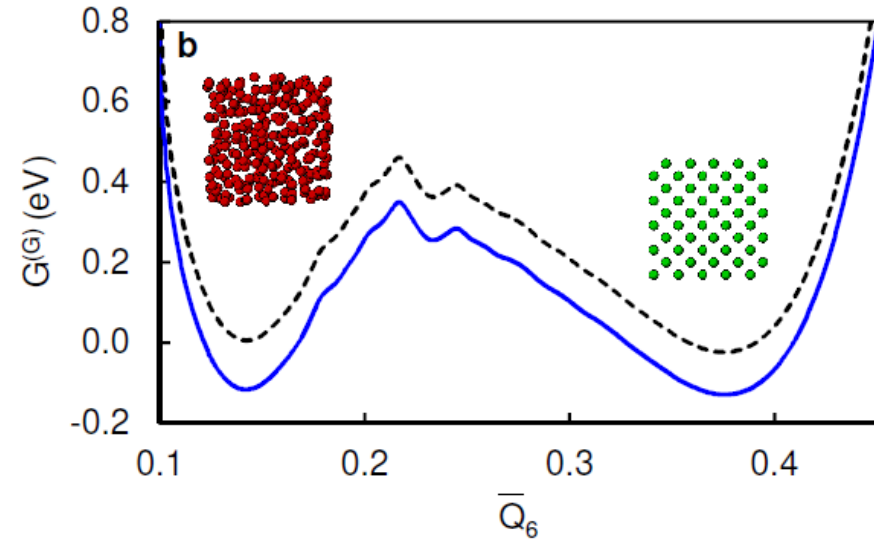
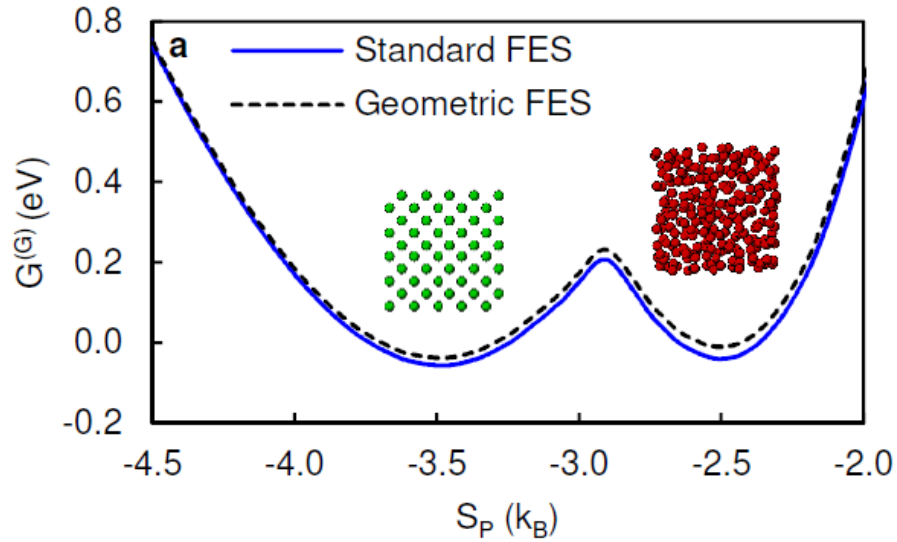


**Encore**

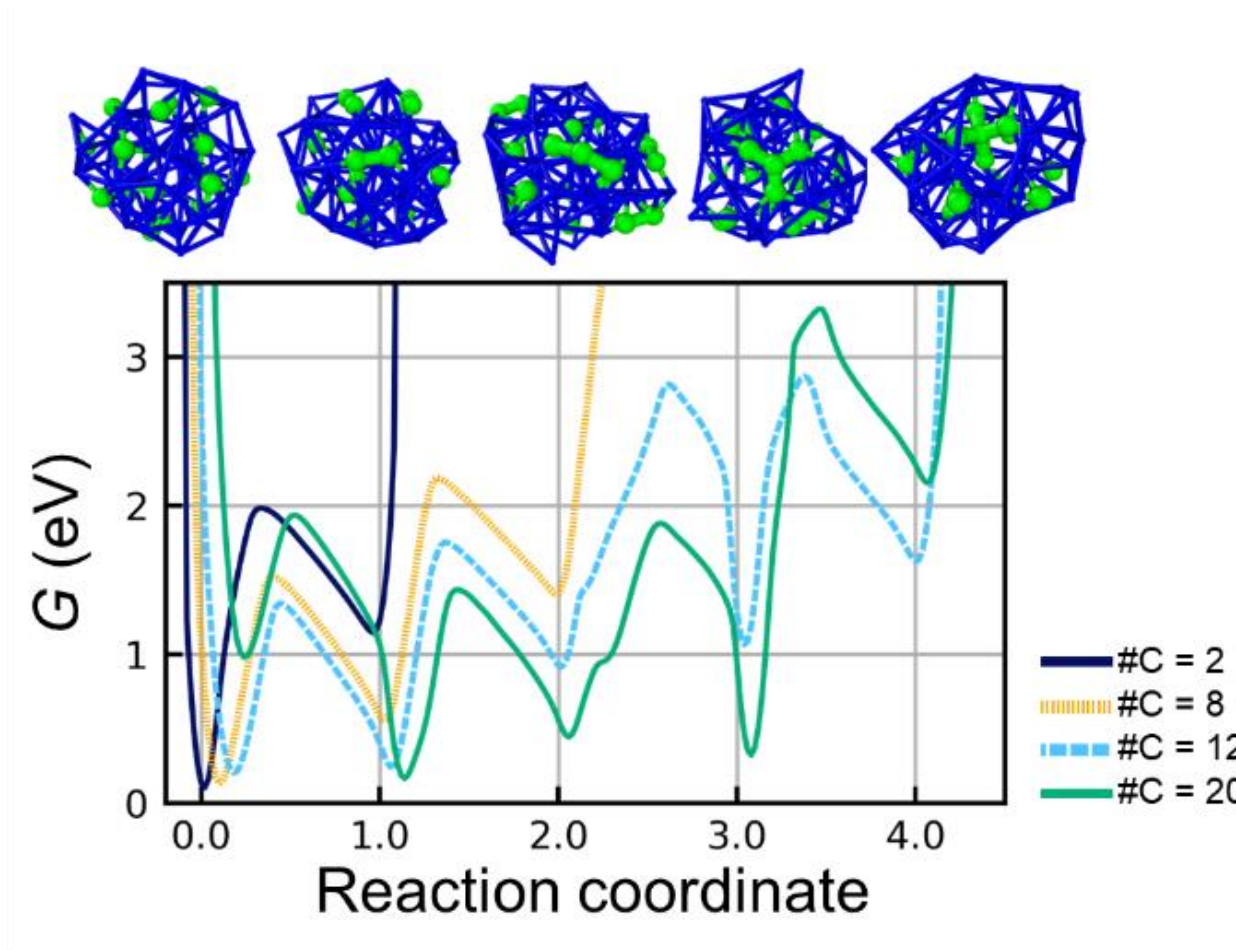
Some extra examples



# Another manybody CV example: liquid Na crystallization



# Initial stages of carbon nanotube nucleation (S. Fukuhara)



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